

Measuring **neutrino mass** imprinted on the **anisotropic galaxy clustering**

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Contents

I. Background Physics

- How **Massive Neutrino** affect **Galaxy Clustering In redshift space?**

II. Theoretical Modeling & Methodology

III. Results

IV. Future work

N_{eff} $\sum m_\nu$

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I. Background Physics

- How **Massive Neutrino** affect **Galaxy Clustering In redshift space?**

Decoupling of Neutrino

$$\rho_R = \sum_{i \in R} g_i \int d^3 p f(p) E(p) = \left[1 + \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} \cdot 3 \right] \rho_\gamma$$

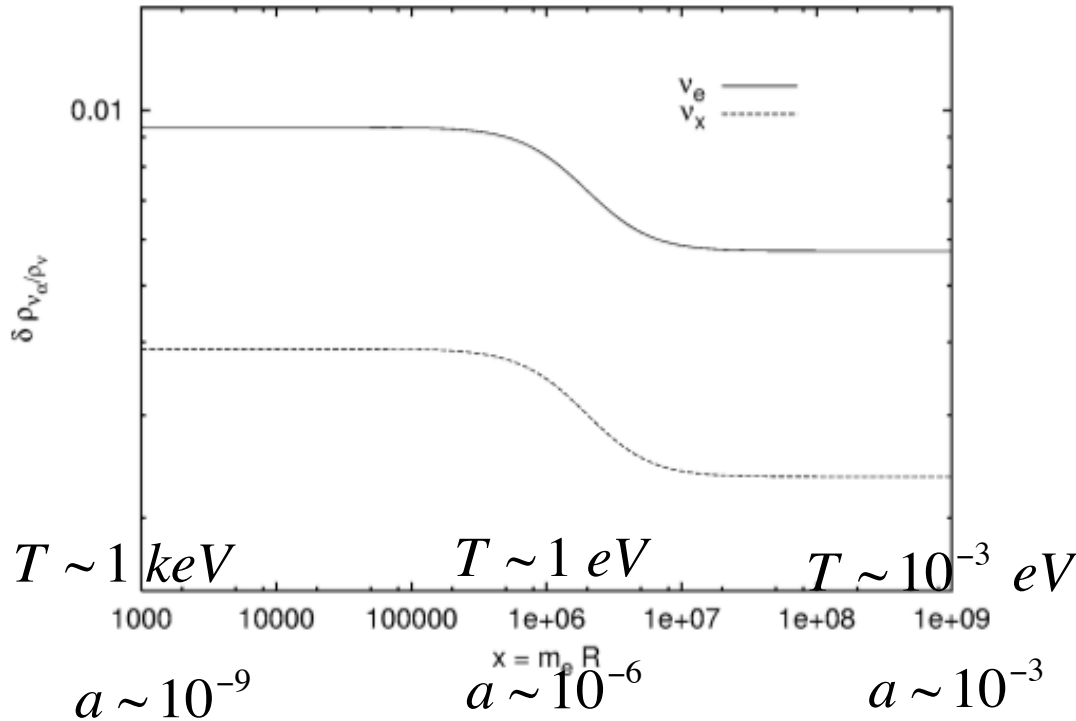
- The contribution from relativistic particles to the energy density.
- If they are in equilibrium with cosmic plasma, FD/BE distribution can be used.
- But, neutrino decoupled at around a few MeV, followed by e-e+ annihilation, which causes heating photons.
- For instantaneous decoupling approximation:

$$T_\gamma / T_\nu = (11/4)^{1/3} \cong 1.40102$$

N_{eff}

Decoupling of Neutrino

$$\rho_R = \sum_{i \in R} g_i \int d^3 p f(p) E(p) = \left[1 + \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} N_{eff} \right] \rho_\gamma$$



$$T_\gamma / T_\nu = (11/4)^{1/3} \cong 1.40102$$

Fig. 1. Evolution of $\delta\rho_{\nu_\alpha}(m_\alpha)/\rho_{\nu_\alpha}^0(m_\alpha)$, for a neutrino mass $m_\alpha = 1$ eV (see text for further details).

Decoupling of Neutrino

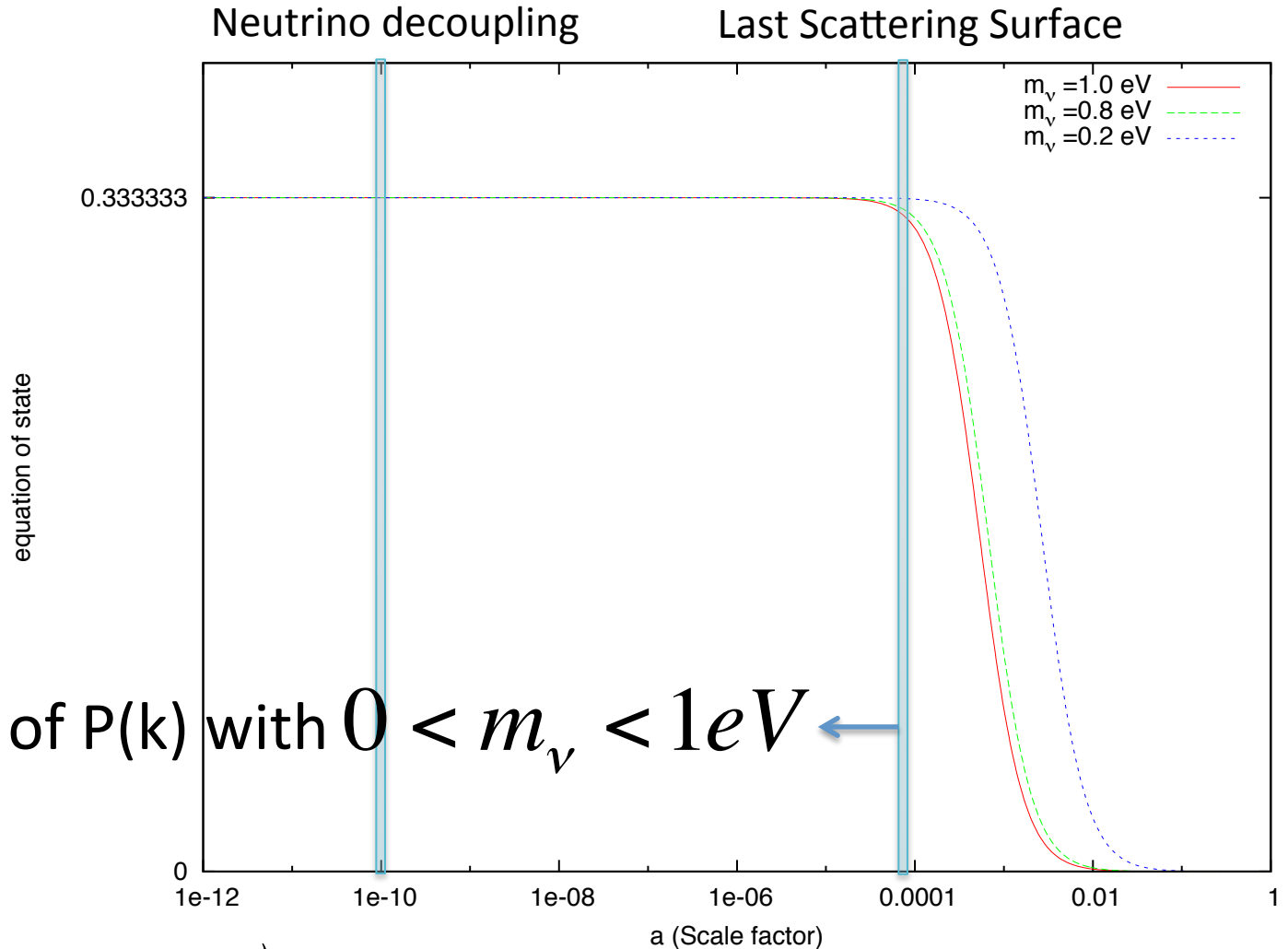
$$\rho_R = \sum_{i \in R} g_i \int d^3 p f(p) E(p) = \left[1 + \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} N_{eff} \right] \rho_\gamma$$

$T_\gamma / T_\nu = (11/4)^{1/3} \cong 1.40102$

- The contribution from relativistic particles to the energy density.
- If they are in equilibrium with cosmic plasma, FD/BE distribution can be used.
- But, neutrino decoupled at around a few MeV, followed by e-e+ annihilation, which causes heating photons.
- Instead of instantaneous decoupling approximation,
- In reality, there is some distortion in f(p),

$$N_{eff} = 3.046$$

After *decoupling*, starts to stream freely



Fix the shape of $P(k)$ with $0 < m_\nu < 1 \text{ eV}$ ←

$$\sum m_\nu \rightarrow m_\nu (\text{degenerate})$$

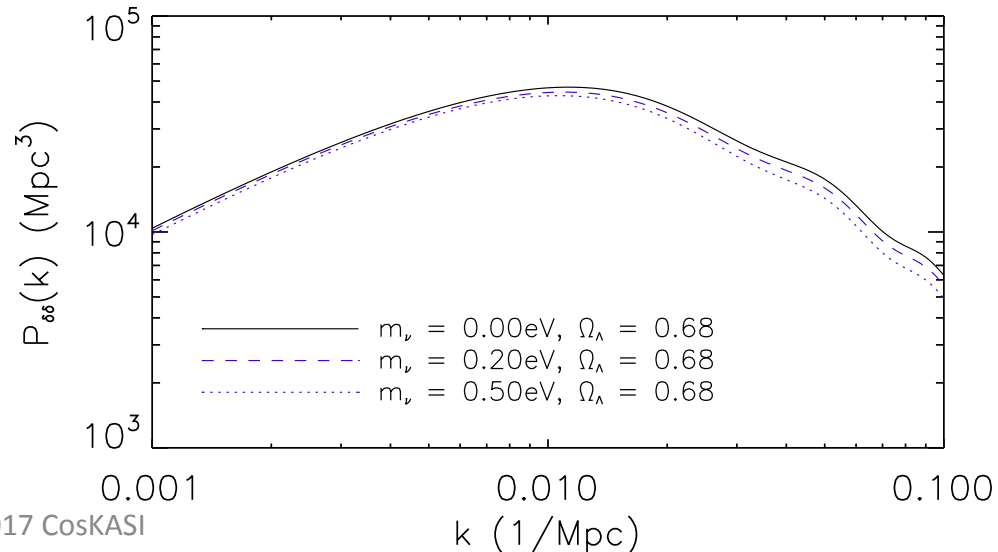
Effect of Massive neutrino on LSS

Yvonne Y. Y. Wong (2011)

- Free-streaming scale of massive neutrino with mass m_ν :

$$\lambda_{FS} \sim 4.2 \sqrt{\frac{1+z}{\Omega_{m,0}}} \left(\frac{eV}{m_\nu} \right) Mpc / h$$

- Structure formation smaller than this scale $k > k_{FS}$ is suppressed, *which provides the access to the neutrino mass in cosmology.*



Contents

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- How **Massive Neutrino** affect **Galaxy Clustering In redshift space**?
- we restrict our analysis to the standard case, where departure of N_{eff} from 3 is solely due to neutrino heating by e-e+ annihilation, which gives The effective number of relativistic species $N_{\text{eff}} = 3.046$.
- Neutrino mass < 1 eV was relativistic before LSS. Therefore, we can fix the clustering feature (=shape of power spectrum) at LSS using Planck experiment result.
- Distortion (scale-dependent damping) from the fixed clustering feature by massive neutrino with $m < 1$ eV provides the access to the neutrino mass in cosmology.

Contents

I. ~~Background Physics~~

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III. Results

IV. Future work

Theoretical model on P(k) in Redshift Space Distortion(RSD)

- Improvement in
2D Power spectrum in redshift space

– Kaiser(1987)

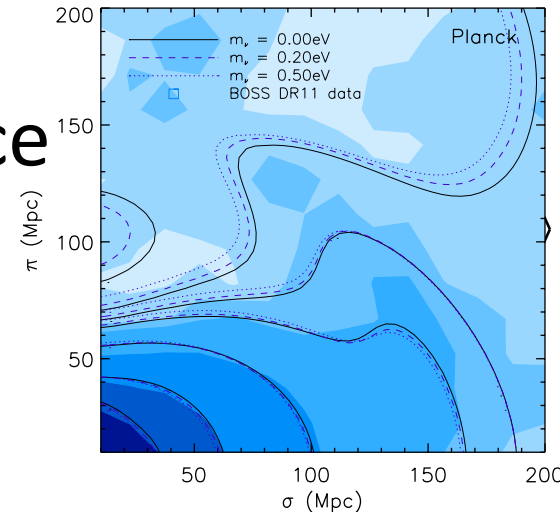
$$P_{Kaiser}^{(s)}(k, \mu) = P_{\delta\delta}^{lin}(k) + 2\mu^2 P_{\delta\Theta}^{lin}(k) + \mu^4 P_{\Theta\Theta}^{lin}(k)$$

– Scoccimarro(2004)

$$P_{scoccimarro}^{(s)}(k, \mu) = \left\{ P_{\delta\delta}(k) + 2\mu^2 P_{\delta\Theta}(k) + \mu^4 P_{\Theta\Theta}(k) \right\} G^{FoG}(k\mu\sigma_p)$$

– Taruya, Nishimichi, and Saito (Improved)(2010)

$$P_{TNS}^{(s)}(k, \mu) = \left\{ P_{\delta\delta}(k) + 2\mu^2 P_{\delta\Theta}(k) + \mu^4 P_{\Theta\Theta}(k) + \underbrace{A(k, \mu) + B(k, \mu)}_{\rightarrow \text{Higher order correction}} \right\} G^{FoG}(k\mu\sigma_p)$$



Theoretical model on P(k) in Redshift Space Distortion(RSD)

$$P^{(s)}(k, \mu) = \int d^3x e^{i\vec{k}\cdot\vec{x}} \langle e^{j_1 A_1} A_2 A_3 \rangle$$

Taruya, Nishimichi, and Saito (2010) *where*

$$j_1 = -ik\mu$$

$$A_1 = u_z(\vec{r}) - u_z(\vec{r}')$$

$$A_2 = \delta(\vec{r}) + \nabla_z u_z(\vec{r})$$

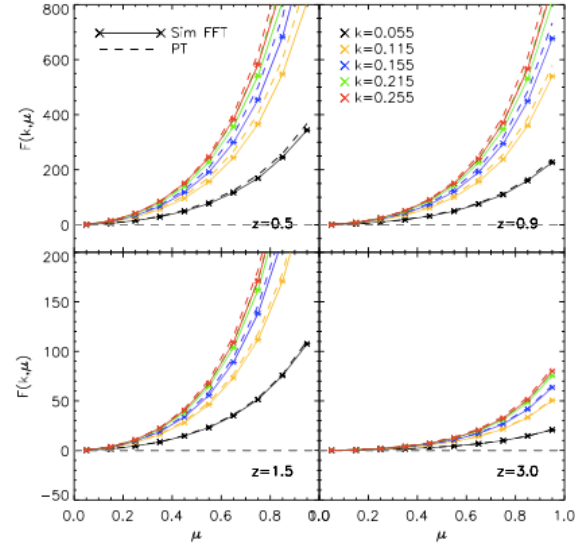
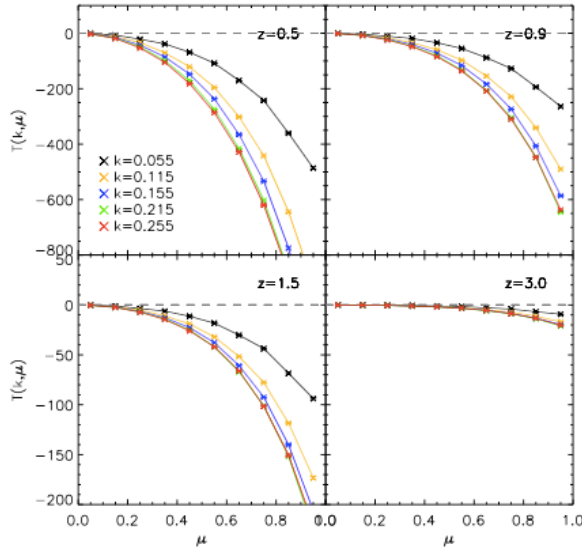
$$A_3 = \delta(\vec{r}') + \nabla_z u_z(\vec{r}')$$

Zheng & Song (2016)

$$P^{(s)}(k, \mu) = \int d^3x \exp\left\{e^{i\vec{k}\cdot\vec{x}} \langle e^{j_1 A_1} \rangle_c\right\} \times \left[\langle e^{j_1 A_1} A_2 A_3 \rangle_c + \langle e^{j_1 A_1} A_2 \rangle_c \langle e^{j_1 A_1} A_3 \rangle_c \right]$$

$$= G^{FoG}(k\mu\sigma_p) \left\{ \begin{aligned} &P_{\delta\delta}(k) + 2\mu^2 P_{\delta\Theta}(k) + \mu^4 P_{\Theta\Theta}(k) \\ &+ A(k, \mu) + B(k, \mu) + T(k, \mu) + F(k, \mu) \end{aligned} \right\}$$

Theoretical model on P(k) in Redshift Space Distortion(RSD)

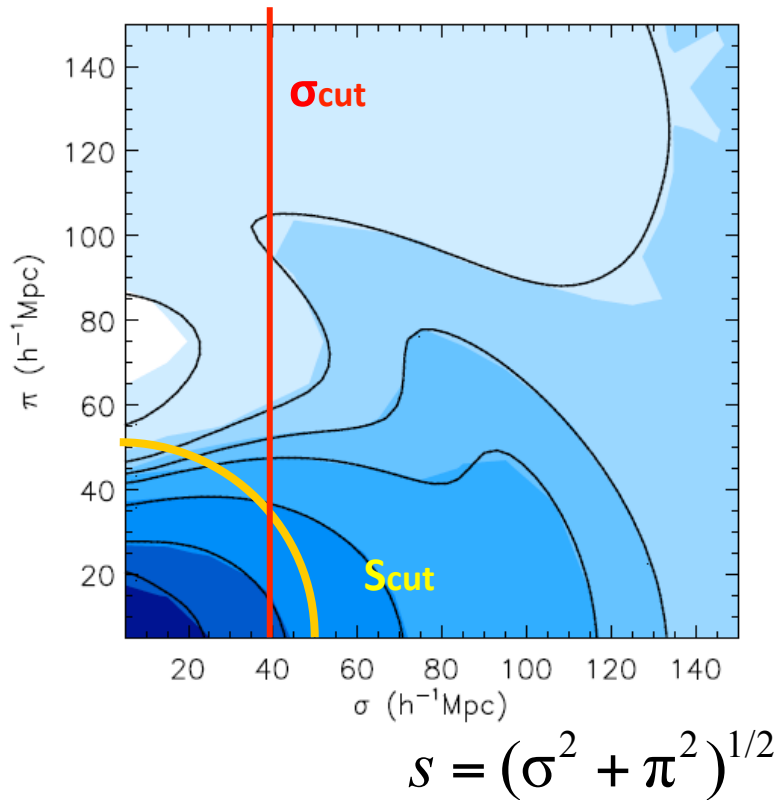


Zheng & Song (2016)

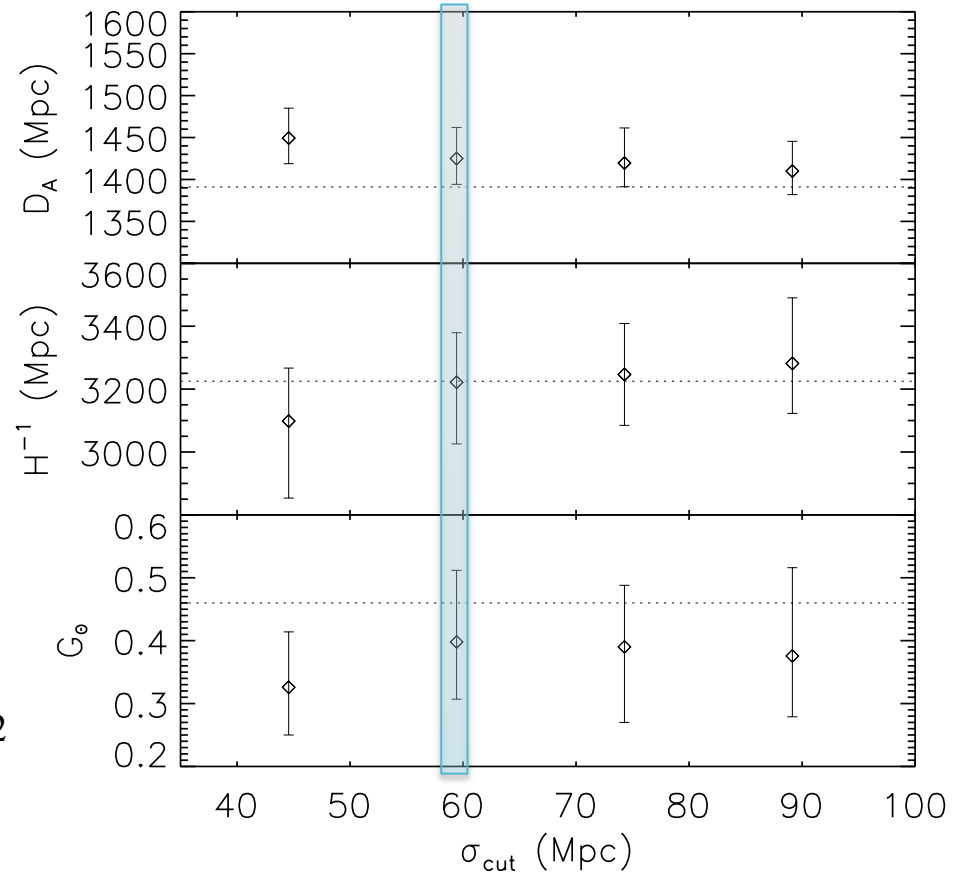
$$\begin{aligned}
 P^{(s)}(k, \mu) &= \int d^3x \exp\left\{e^{i\vec{k}\cdot\vec{x}} \left\langle e^{j_1 A_1} \right\rangle_c\right\} \times \left[\left\langle e^{j_1 A_1} A_2 A_3 \right\rangle_c + \left\langle e^{j_1 A_1} A_2 \right\rangle_c \left\langle e^{j_1 A_1} A_3 \right\rangle_c \right] \\
 &= G^{FoG}(k\mu\sigma_p) \left\{ \begin{aligned} &P_{\delta\delta}(k) + 2\mu^2 P_{\delta\Theta}(k) + \mu^4 P_{\Theta\Theta}(k) \\ &+ A(k, \mu) + B(k, \mu) + T(k, \mu) + F(k, \mu) \end{aligned} \right\}
 \end{aligned}$$

Cut-off to consider

Current status of RSD modeling



$$k_{\text{max}} \sim 0.07h / \text{Mpc}$$

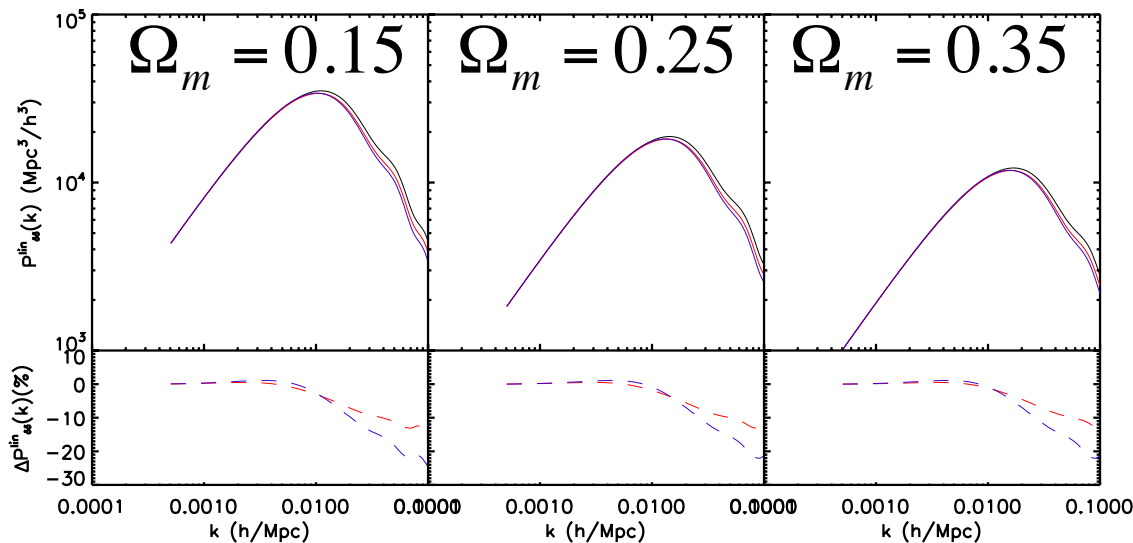


Is TNS model **reliable** to calculate Non-linear mapping including **massive neutrino**?

$$\frac{\Delta P}{P} \sim 8 \times f_\nu = 0.16 \text{ for } m_\nu = 0.3 \text{ eV}$$

$$\text{where } f_\nu \equiv \frac{\Omega_\nu}{\Omega_m} \text{ with } \Omega_m = 0.31 \text{ and } h = 0.68$$

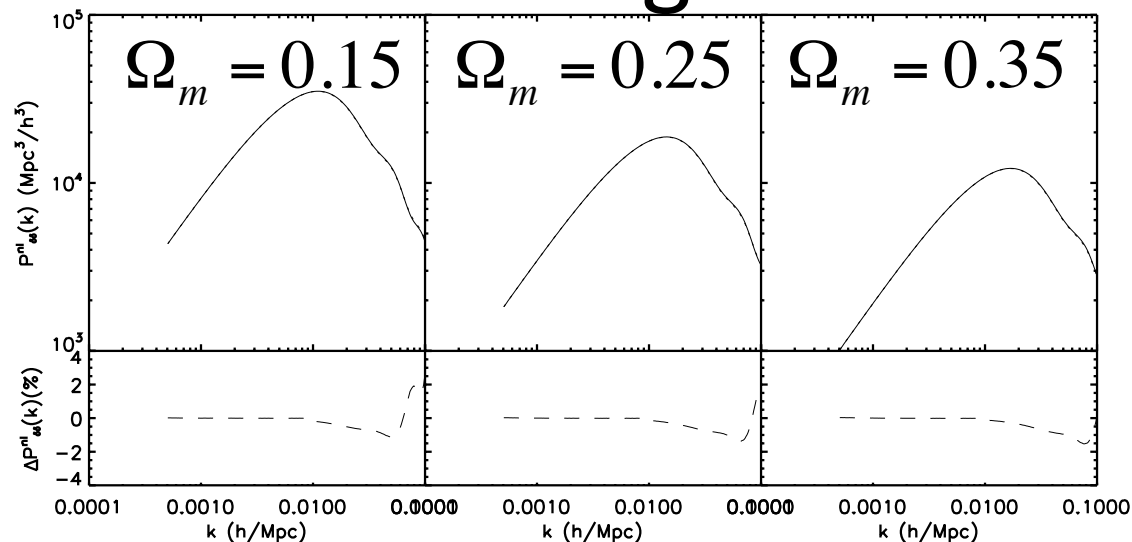
Yvonne Y. Y. Wong (2011)



- $m_\nu = 0.0$ eV
- $m_\nu = 0.3$ eV ←
- $m_\nu = 0.6$ eV

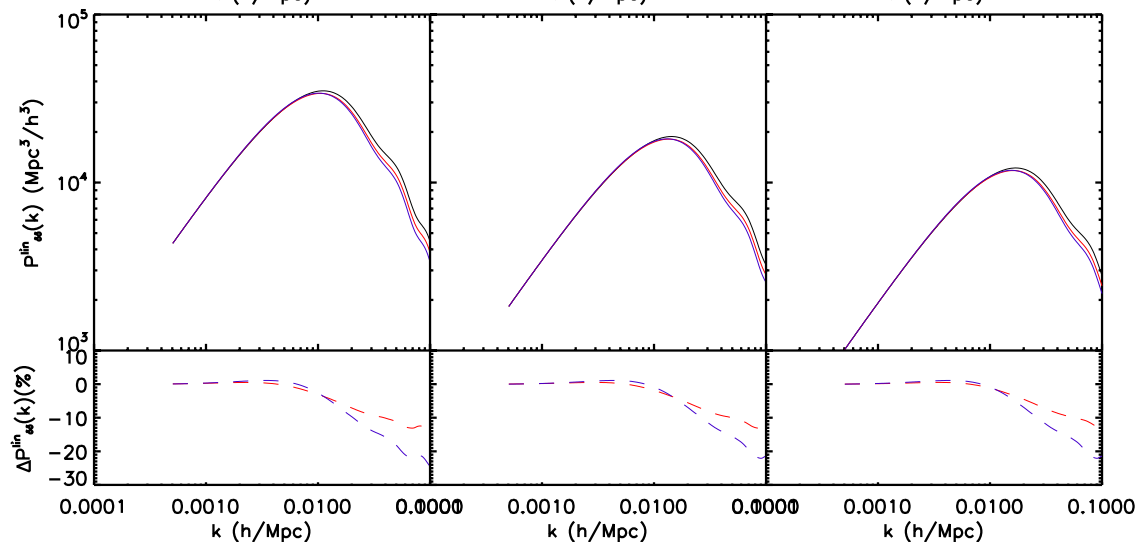
Fractional difference between linear $P_{\delta\delta}(k)$ without and with massive neutrino $\sim 15\%$.

Is TNS model **reliable** to calculate Non-linear mapping including **massive neutrino**?



→ Yes
up to $k \sim 0.1$ h/Mpc

Fractional difference between linear and non-linear $P_{\delta\delta}(k)$ without massive neutrino $< 5\%$.

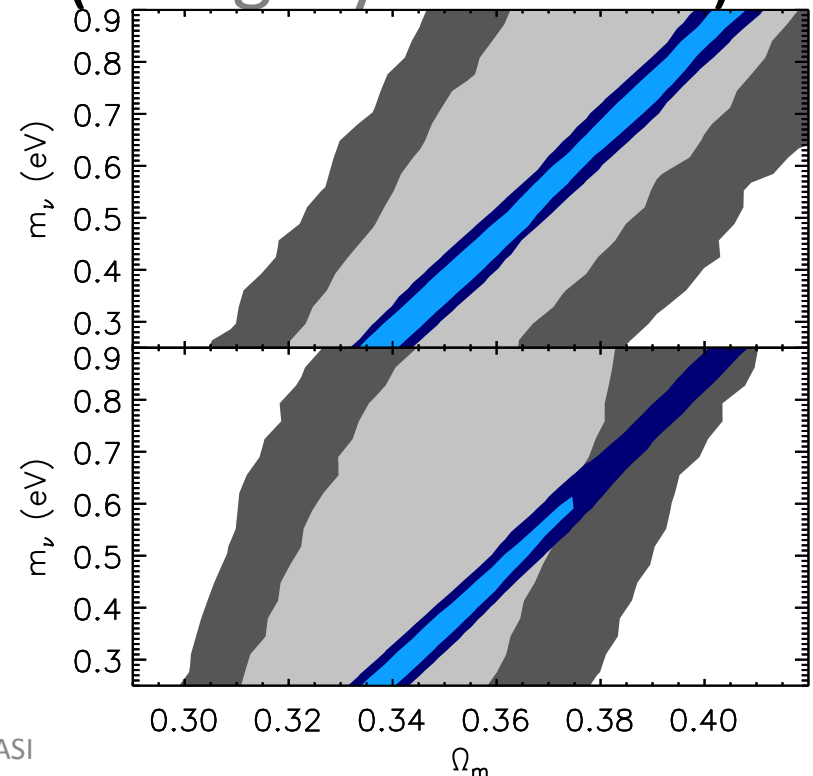


- $m_\nu = 0.0$ eV
- $m_\nu = 0.3$ eV
- $m_\nu = 0.6$ eV

Fractional difference between linear $P_{\delta\delta}(k)$ without and with massive neutrino $\sim 15\%$.

Effective growth VS Scale-dep. growth

- Depending on how the effect from massive neutrino is parameterized, the constraint on neutrino mass is affected (See grey contours).



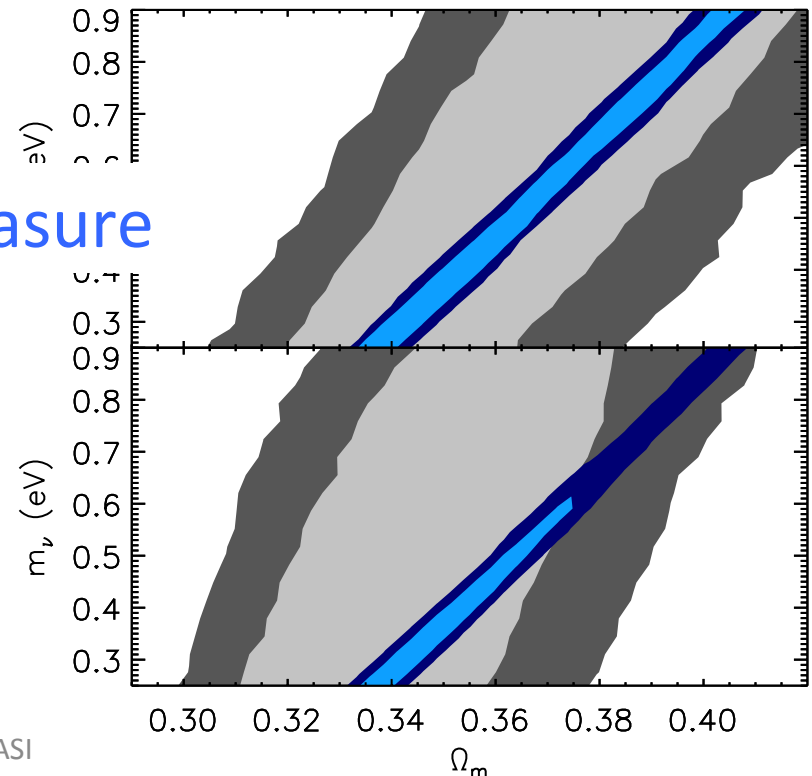
Effective growth VS Scale-dep. growth

- Depending on how the effect from massive neutrino is parameterized, the constraint on neutrino mass is affected.

$$\theta_* = \frac{r_s(z_*)}{D_A(z_*)} : \text{CMB distance measure}$$

$$D_A(a) = a \int_a^1 \frac{da'}{a' H(a')}$$

$$r_s(z) = \int_0^{\eta(z)} \frac{d\eta'}{\sqrt{3(1+R)}}$$



Bias effect on neutrino mass constraint

- Beyond the linear bias, b_1 ?

$$P_{g,\delta\delta}(k) = b_1^2 P_{\delta\delta}(k) + 2b_2 b_1 P_{b_2,\delta}(k) + b_2^2 P_{b_22}(k) + 2b_{s_2} b_1 P_{b_{s_2},\delta}(k) + 2b_2 b_{s_2} P_{b_2 s_2}(k) + b_{s_2}^2 P_{b_{s_2}2}(k) + 2b_{s_2} b_{3nl} \sigma_3^2(k) P^{lin}(k)$$

b1 and b2: local bias

b_{s2}, and b_{3nl}: non-local bias

where $P_{b_2,\delta}(k) = \int \frac{d^3 q}{(2\pi)^3} P^{lin}(q) P^{lin}(|\vec{k} - \vec{q}|) F_2^{SPT}(\vec{q}, \vec{k} - \vec{q})$

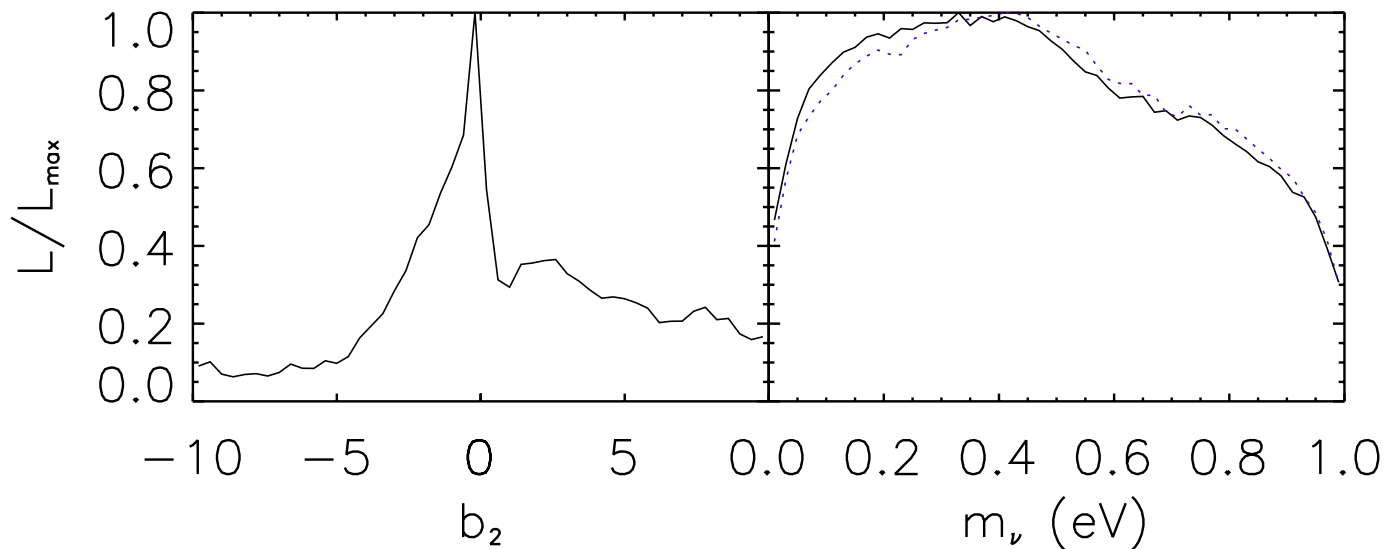
$$P_{b_22}(k) = -\frac{1}{2} \int \frac{d^3 q}{(2\pi)^3} P^{lin}(q) \left[P^{lin}(q) - P^{lin}(|\vec{k} - \vec{q}|) \right]$$

McDonald & Roy (2009)
Gill-Marín+ (2016)

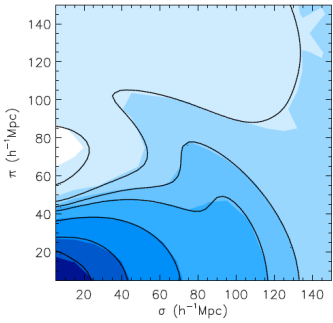
Bias effect on neutrino mass constraint

- Scale-dependency of bias $b(k)$ doesn't affect neutrino mass constraint in scales of interest.

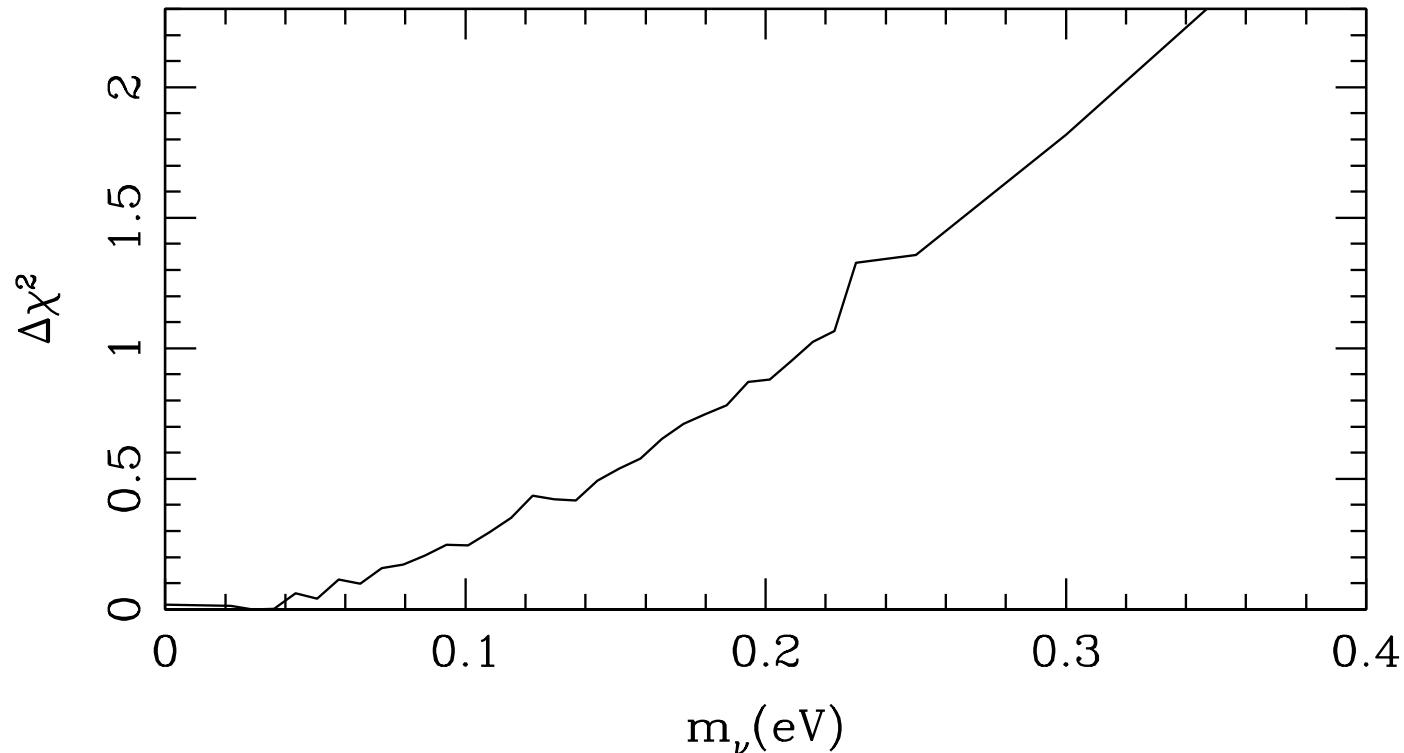
$$P_{g,\delta\delta}(k) = b_1^2 P_{\delta\delta}(k) + 2b_2 b_1 P_{b2,\delta}(k) + b_2^2 P_{b22}(k)$$



Testing Methodology



- When we apply our methodology to the simulation (SDSS DR11 mock catalogue without massive neutrino), true value reproduced.



Contents

I. ~~Background Physics~~

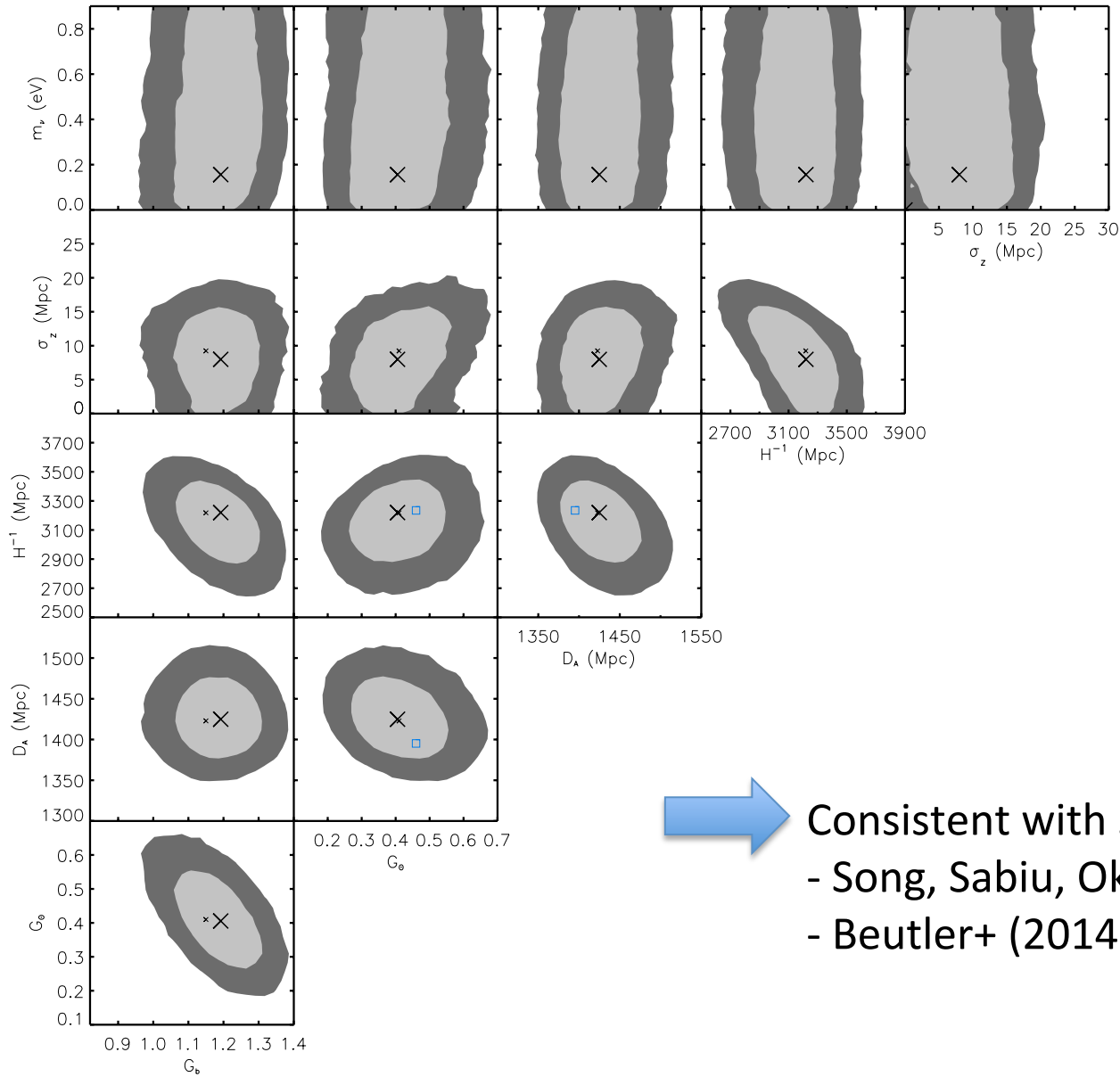
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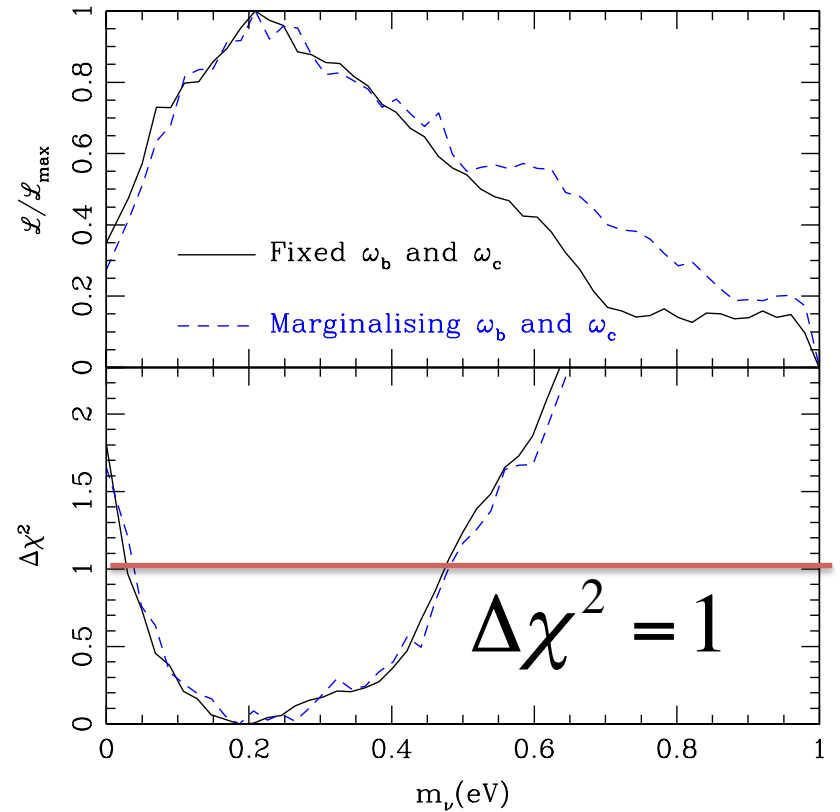
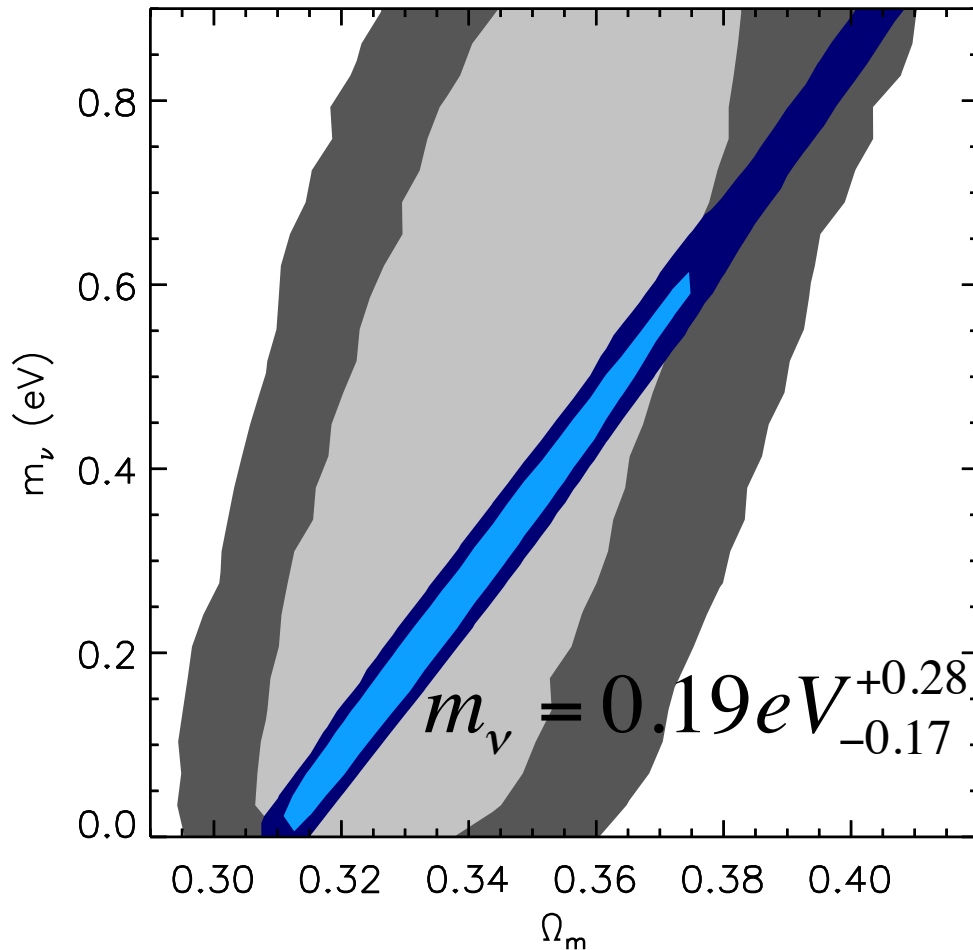
Result on $(D_A, H^{-1}, G_b, G_\theta, \sigma_p, m_\nu)$



Consistent with some previous works:

- Song, Sabiu, Okumura, Oh, Linder (2014)
- Beutler+ (2014)

Result on $(b, \Omega_m, \sigma_p, m_\nu) + \theta_*$ in 68% C. L.



$$m_\nu = 0.22 eV^{+0.28}_{-0.17} \text{ (HPD)}$$

$$m_\nu = 0.31 eV^{+0.16}_{-0.26} \text{ (equal-tailed)}$$

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Future Work

- Precision in Theoretical prediction
 - To prepare forthcoming **DESI** data with higher precision, theoretical prediction for nonlinearity in redshift space should be more elaborated up to higher k where the effect of massive neutrino comes in.
 - Alternatively, templates could be supplied by neutrino simulations. (similar manner to Zheng & Song 2016)

Future Work

- Precision in Theoretical prediction
- with full-scale information from CMB instead of one distance scale
- using SDSS DR12.

Summary

- The effect of massive neutrino with mass < 1 eV, which decoupled when it was relativistic & became non-relativistic after LSS, affect anisotropic galaxy clustering (SDSS DR11 CMASS at $z_{\text{eff}} = 0.57$), which let us access neutrino mass to give $m_\nu = 0.19 eV_{-0.17}^{+0.28}$ in 68% C.L.
 - TNS model is available for massive neutrino with $k_{\text{max}} < 0.1$.
 - Our results are conservative in change of local bias.
 - Free form of Dark energy doesn't help us to constrain neutrino mass, but consistent with the previous works.
 - Cosmological constant with CMB distance measure can help us for neutrino mass.
 - Type of credible/confidence Interval doesn't change much reporting our results.

Thank you!