# Gravity from Cosmology 

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## The state of General Relativity in 1957

"There exists... one serious difficulty, and that is the lack of experiments. Furthermore, we are not going to get any experiments, so we have to take the viewpoint of how to deal with the problems where no experiments are available. ... the best viewpoint is to pretend that there are experiments and calculate. In this field we are not pushed by experiments but pulled by imagination."



"The elegant logic of general relativity theory, and its precision tests, recommend GR as the first choice for a working model for cosmology. But the Hubble length is fifteen orders of magnitude larger than the length scale of the precision tests, at the astronomical unit and smaller, a spectacular extrapolation."

Jim Peebles, IAU 2000


(Baker, Psaltis \& Skordis 2014)

## Big puzzles ...



# small inconsistencies. 

- Lamb shift
- Wu parity violation experiment
- Fitch-Cronin CP violation
- Precession of perihelion of Mercury


## The Theory

## Einstein Gravity

$$
\frac{\text { Curvature }}{16 \pi G} \int d^{4} x \sqrt{-g}[R(g)-2 \Lambda]+\int d^{4} x \sqrt{-g} \mathcal{L}(g, \text { matter })
$$

Lovelock's theorem (197I):"The only second-order, local gravitational field equations derivable from an action containing solely the 4D metric tensor (plus related tensors) are the Einstein field equations with a cosmological constant."

See also Hojman, Kuchar \& Teitelboim (1976)

## Jordan-Brans-Dicke Theory

## One free parameter



Cassini (Bertotie ea a 2003) $\omega_{\mathrm{BD}}>40,000$
Einstein-Dilaton-Gauss-Bonnet
Cascading gravity

Randall-Sundrum I \& II

Higher dimensions


Generalisations

Gauss-Bonnet

Lovelock gravity

## arXiv:

1310.1086
1209.2117

I 107.0491
| | 10.3830
of $S_{\text {EH }}$
Teves - New degrees of freedom

Some degravitation scenarios

Conformal gravity

$$
f(G)
$$

General $R_{\mu \nu} R^{\mu \nu}$, $\square R$,etc.

## Extra degrees of freedom

 metric $\longrightarrow$ add $\phi, A^{\mu}, f_{\alpha \beta}$ etc.$4 \mathrm{D} \longrightarrow$ e.g. in 5 dimensions:

$$
g_{A B}=\left(\begin{array}{cc}
g_{\alpha \beta} & A_{\alpha} \\
A_{\beta} & \phi
\end{array}\right)
$$

2nd order $\longrightarrow$ e.g. if $\int d^{4} x \sqrt{-g} f(R)$ define $\phi=\frac{d f}{d R}$.


Background $d s^{2}=g_{\mu \nu}^{(0)} d x^{\mu} d x^{\nu}=-d t^{2}+a^{2}(t)(d \vec{x})^{2}$

$$
G_{\alpha \beta}=8 \pi G T_{\alpha \beta}+\underline{U_{\alpha \beta}}
$$

where $U_{\alpha \beta}\left(a, \dot{a}, \rho_{M}, P_{M}, \phi, \cdots\right) \quad P_{X}=\underline{w} \rho_{X}$
Homogeneity and isotropy
BOSS, Anderson et al 2013.

$$
U_{\alpha \beta}=8 \pi G\left(\begin{array}{cc}
\rho_{X} & 0 \\
0 & a^{2} P_{X} \delta_{i j}
\end{array}\right)
$$

Bianchi identities

$$
\nabla^{\alpha}\left[8 \pi G T_{\alpha \beta}+U_{\alpha \beta}\right]=0
$$


$\frac{\delta T}{T} \sim 10^{-5}$

## linear perturbation theory

## Linear Perturbations

$$
\begin{aligned}
g_{\alpha \beta} & =g_{\alpha \beta}^{(0)}+h_{\alpha \beta} \\
\rho_{M} & =\bar{\rho}_{M}\left(1+\delta_{M}\right) \\
\phi & =\phi_{0}+\delta \phi
\end{aligned}
$$

Construct most general quadratic action which has:

- upto 2 nd order in time derivatives
- $h_{\alpha \beta} \rightarrow h_{\alpha \beta}+\nabla_{\alpha} \xi_{\beta}+\nabla_{\beta} \xi_{\alpha}$ where $x^{\alpha} \rightarrow x^{\alpha}+\xi^{\alpha}$
- inherits symmetries of the background


## Linear Perturbations

$$
\begin{aligned}
& S=\int d^{4} x \sqrt{-g^{(0)}}\left[\alpha_{1} \nabla^{\mu} h^{\alpha \beta} \nabla_{\mu} h_{\alpha \beta}+\alpha_{2} \nabla^{\mu} h^{\alpha \beta} \nabla_{\alpha} h_{\mu \beta}\right. \\
& \left.+\alpha_{3} \nabla^{\mu} h \nabla^{\alpha} h_{\mu \alpha}+\alpha_{4} \nabla^{\mu} h \nabla_{\mu} h+\alpha_{5} \nabla^{\mu} h_{\mu \alpha} \nabla^{\alpha} \delta \phi+\cdots\right]
\end{aligned}
$$

## Properties:

- $\left(\alpha_{1}, \alpha_{2}, \cdots\right)$ are functions of $t$
- $\alpha_{X}(t)$ depend on transf. props of extra fields
- clear mapping theory $\longleftrightarrow \alpha_{X}(t)$
- clear physical interpretation of each $\alpha_{X}(t)$

Examples:

- Scalar-tensor (Horndeski): five $\alpha_{X}(t)$
- Vector-tensor (Einstein-Aether, Proca): nine $\alpha_{X}(t)$
- Tensor-tensor (Bigravity, massive gravity): three $\alpha_{X}(t)$


$N(k) \propto k^{3}$
More statistical power


## The Data

A preferred length scale- the horizon

$$
\mathcal{H}^{-1} \equiv\left(\frac{\dot{a}}{a}\right)^{-1} \propto \tau \simeq 3000 h^{-1} \mathrm{Mpc}
$$

Most surveys $\ll \tau$ so that $k \tau \gg 1$
Newtonian potentials: $h_{\alpha \beta}=2\left(\begin{array}{cc}\Phi & 0 \\ 0 & a^{2} \Psi \delta_{i j}\end{array}\right)$
Einstein equations: $-k^{2} \Phi=4 \pi G \underline{\mu} a^{2} \rho \Delta$

$$
\underline{\gamma} \Psi=\Phi
$$

$(\mu, \gamma)$ are rational functions of $\alpha_{X}(t)$ and $k^{2}$

## We measure matter and light.



## Growth rate

$$
f(k, a)=\frac{d \ln \delta_{M}(k, a)}{d \ln a}
$$

$f$ satisfies a simple ODE


$$
\frac{d f}{d \ln a}+q f+f^{2}=\frac{3}{2} \Omega_{M} \xi
$$

with $q=\frac{1}{2}\left[1-3 w\left(1-\Omega_{M}\right)\right]$ and $\xi=\frac{\mu}{\gamma}$


## Weak Lensing



$$
\begin{aligned}
\text { shear } & \simeq \int_{0}^{\chi} \nabla_{\perp}^{2}[\Phi+\Psi]\left(\chi^{\prime}\right)\left[\chi^{\prime}\left(1-\frac{\chi^{\prime}}{\chi}\right)\right] d \chi^{\prime} \\
\text { shear } & \sim \Sigma \equiv \mu\left(1+\frac{1}{\gamma}\right)
\end{aligned}
$$

Sarah Bridle lectures (2003)

Joudaki et al 2016


## Lensing of CMB



Planck 2015


Constrain $\alpha_{X}(t)$ in scalar-tensor theories
Parametrize: $\alpha_{X}=b_{X}+c_{X} \frac{\Omega_{\mathrm{DE}}(z)}{\Omega_{\mathrm{DE}}(z=0)}$



$$
\sigma\left(\alpha_{X}\right) \sim 0.5
$$



Bellini et al 2016

## Jordan-Brans-Dicke Theory

## One free parameter

$S=\int d^{4} x \sqrt{-g}\left[\phi R-\frac{\omega_{\mathrm{BD}}}{\phi} \nabla_{\mu} \phi \nabla^{\mu} \phi+V+\mathscr{L}_{M}\left[g_{\mu \nu}, \varphi\right]\right\}$

Cassini ${ }_{\text {(Bertotti et al 2003) }} \omega_{\mathrm{BD}}>40,000$
Planck (Avilez \& Skordis 2015) $\omega_{\mathrm{BD}}>1,000$

## The Challenge

## Systematics: non-linear physics


$N(k) \propto k^{3}$
More statistical power

## Systematics: non-linear physics

baryonic feedback


Sembloni et al 2012

## Systematics: non-linear physics



Jennings, Baugh \& Pascoli 2015

## Systematics: screening

Newtonian potential
GM
$r$

Fifth force
$\phi=-\frac{\tilde{G} M}{r} e^{-m r}$
Chameleon: $m=m(\rho)$
$m \rightarrow \infty$ when $\rho \gg \rho_{S}$
Vainshtein: $\quad \tilde{G} \rightarrow 0$ when $r \ll r_{V}$

## The Future

## The Future is now

| Data Type | Now | Soon | Future |
| :---: | :---: | :---: | :---: |
| Photo-z:LSS <br> (weak lensing) | DES, RCS, KIDS | HSC | LSST, Euclid, SKA, <br> WFIRST |
| Spectro-z <br> (BAO, RSD, ...) | BOSS | DESI,PFS,HETDEX, <br> Weave | Euclid, SKA |
| SN la | HST, Pan-STARRS, <br> SCP, SDSS, SNLS | DES, J-PAS | JWST,LSST |
| CMB/ISW | WMAP, Planck | AdvACT | Simons Array, Stage <br> IV, LiteBird |
| sub-mm, small scale <br> lensing, SZ | ACT, SPT,Planck, <br> ACTPol,SPTPol, | PolarBear,Spider, <br> Vista | CCAT, SKA |
| X-Ray clusters | ROSAT, XMM, <br> Chandra | XMM, XCS, eRosita |  |
| HITomography | GBT | Meerkat, Baobab, <br> Chime, Kat 7 | SKA |



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"An important contribution of the general theory of relativity to cosmology has been to keep out theologians by a straightforward application of tensor analysis."

E. Schucking

D. Alonso, L.Amendola, M.Amin,T. Baker, R.Bean, E. Bellini, C. Blake, P. Bull, P. Brax, S. Daniels, A. Davies, D. Leonard, G. Gubitosi, P. G. Ferreira, J. Gleyzes,W. Hu, L. Hui, C, Heymans, S. Joudaki, K. Koyama, M. Kunz, M. Lagos, D. Langlois, E. Linder, L. Lombrisier, D. Mota, A. Narimani, J. Noller, J. Peacock, F. Piazza, D. Pogosian, D. Sapone, D. Scott, I. Sawicki, A. Silvestri, F. Simpson, A. Taylor, F.Vernizzi, H.Winther, J. Zuntz, ...

