Nonlinear velocity statistics and redshift-space distortions in peculiar velocity surveys

CosKASI Conference 2014 Apr. 16-18, 2014

Teppei OKUMURA

Kavli Institute for the Physics and Mathematics of the Universe (Kavli iPMU), The University of Tokyo

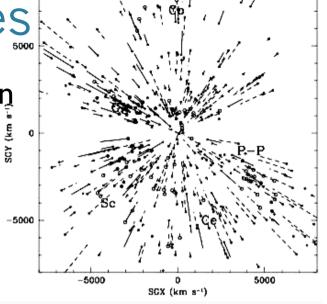
Collaborators: Uros Seljak (Berkeley), Zvonimir Vlah (Zurich), Vincent Desjacques (Geneva)

Okumura, Seljak, Vlah, Desjacques, (2014) JCAP accepted (arXiv:1312.4214)

Peculiar velocities of galaxies

Contains information of 3D mass distribution

Distance to galaxies $z = H_0 x$ Radial peculiar through spectroscopy Redshift True distance (redshift space) (real space)

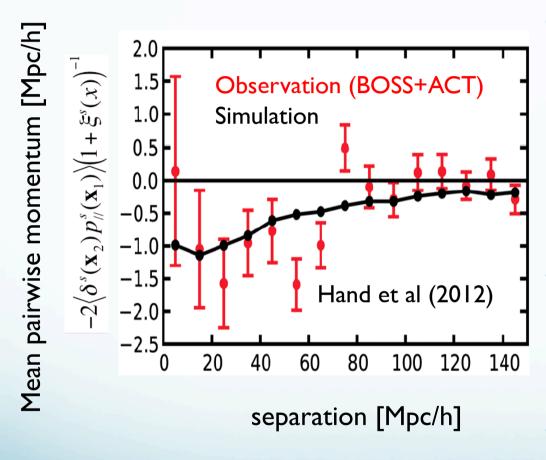


 It's hard to measure peculiar velocities themselves, but by measuring power spectrum in galaxy surveys the effect of peculiar velocities shows up as anisotropy (Redshift-sapce distortions)

Formula of galaxy power spectrum in $P_g^s(k,\mu) = (b + f\mu^2)^2 P_m^r(k)$ Innear thoery (Kaiser) Anisotropy due to RSD (μ : direction cosine between wavevector and LoS)

Theoretical modeling of redshift-space density power spectrum has been extensively studies as an observable.

Direct measurement of peculiar velocities: kinetic Sunyaev-Zeld'vich (kSZ)

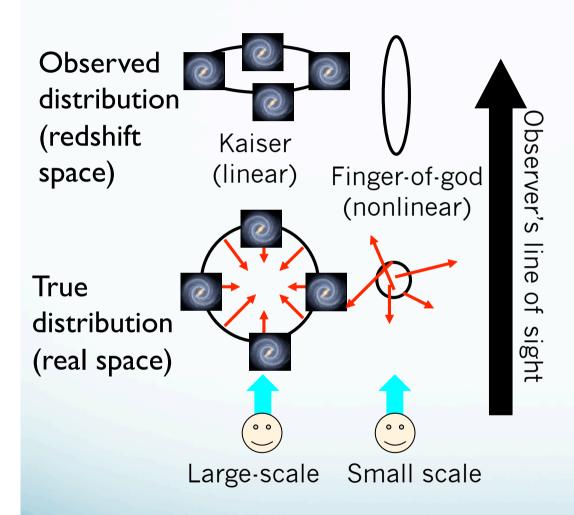


- By combining a galaxy survey (BOSS) and a CMB map (ACT), measured the distortion of the CMB spectrum along the line of observed clusters of galaxies (kSZ effect) and determined peculiar velociteis
- Observables are momentum field sampled in redshift-space positions of galaxies, thus affected by (nonlinear) RSD.
- However, such a formulation has been made so far only by linear theory in Fourier space (thus simply Kaiser).

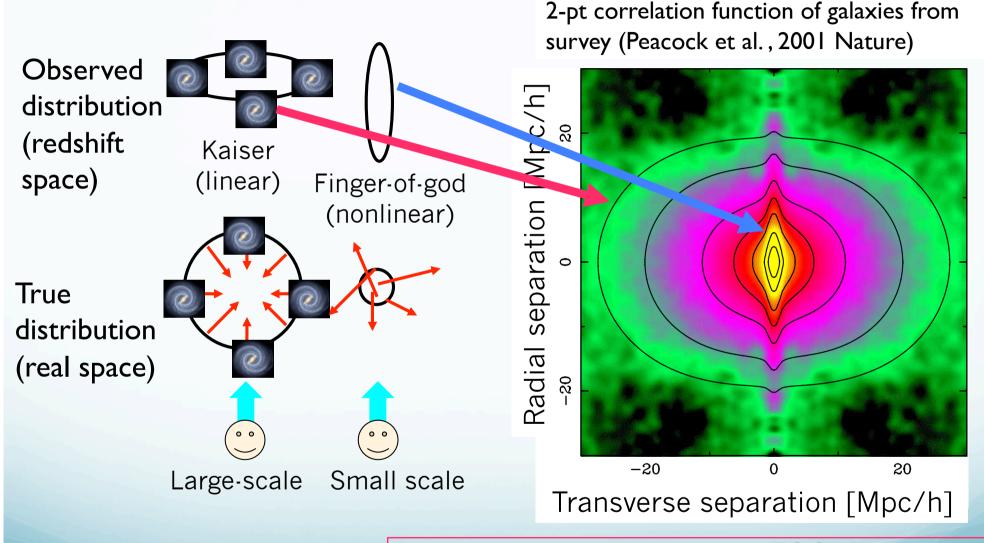
Purpose and conclusion of this study

- We define the momentum field in redshift space, and derive formulae for peculiar velocity statistics that include nonlinear RSD.
- We compute the statistics using Eulerinan 1-loop perturbation theory and linear theory.
- We compare them to the numerical results measured from N-body simulations.
- We find that N-body results cannot be explained by linear theory even in very large scales, while our nonlinear PT improves the accuracy.
 - reference:
 - Okumura, Seljak, Vlah, Desjacques, (2014) JCAP (arXiv:1312.4214)

Redshift-space distortions (RSD)



Redshift-space distortions (RSD)



Linear power spectrum + FOG damping $P_g^s(k,\mu) = (b + f\mu^2)^2 P_m^r(k) \times \exp(-k^2 \mu^2 \sigma_v^2)$

Redshift-space power spectrum using phase-space distribution function

RSD is defined in phase space $f(\mathbf{x}, \mathbf{p})$ starting from Vlasov equation

$$\vec{s} = \vec{x} + \hat{z}v_{//}/aH = \vec{x} + \hat{z}p_{//}/a^2mH$$
 Distortions due to radial momentum (~10Mpc/h)

momentum (~10Mpc/h)

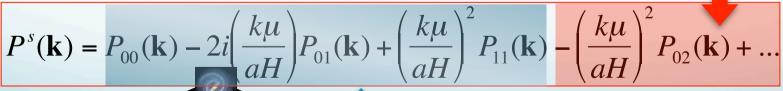
Redshfit-space density field is described by real-space one + infinite sum of moments of mass-weighted velocities

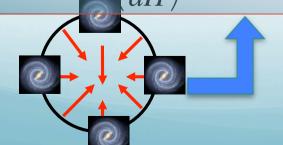
$$\delta^{s}(\mathbf{k}) = \delta^{r}(k) + \sum_{L=1}^{\infty} \frac{1}{L!} \left(\frac{i k_{//}}{a H} \right) T_{//}^{L}(\mathbf{k})$$
 where

$$\delta^{s}(\mathbf{k}) = \delta^{r}(k) + \sum_{L=1}^{\infty} \frac{1}{L!} \left(\frac{i k_{//}}{a H} \right) T_{//}^{L}(\mathbf{k}) \text{ where } T_{//}^{L}(\mathbf{x}) = \frac{m}{\overline{\rho} a^{3}} \int d^{3} p \ f(\mathbf{x}, \mathbf{p}) \left(\frac{p_{//}}{a m} \right)^{L} = \left[1 + \delta(\mathbf{x}) \right] u^{L}(\mathbf{x})$$

Power spectrum becomes an infinite sum as well.

$$\langle \delta^{s}(\mathbf{k})\delta^{s}(\mathbf{k}')\rangle = (2\pi)^{3}P^{s}(\mathbf{k})\delta^{D}(\mathbf{k} - \mathbf{k}')$$





Seljak & McDonald (2011) JCAP (paper I)

Okumura, Seljak, McDonald, Desjacques (2012a) JCAP (II)

Okumura, Seljak, Desjacques (2012,b) JCAP (III)

Vlah, Seljak, McDonald, Okumura, Baldauf (2012) JCAP (IV)

Vlah, Seljak, **Okumura**, Desjacques (2013) JCAP (V)

Extension to redshift-space momentum

$$\vec{s} = \vec{x} + \hat{z}v_{//}/aH = \vec{x} + \hat{z}p_{//}/a^2mH$$
 Distortions due to radial

peculiar velocities (~10Mpc/h)

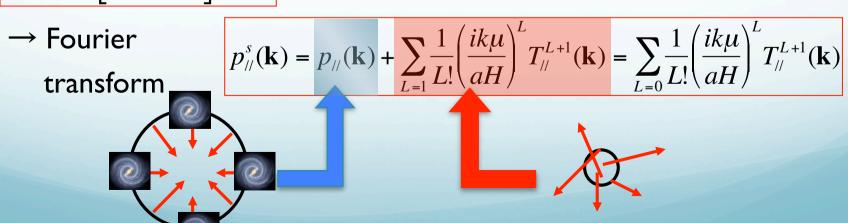
Okumura, Seljak, Vlah, Desjacques (arXiv:1312.4214)

Derivation for density field can be applied to momentum

$$\delta^{s}(\mathbf{k}) = \delta^{r}(k) + \sum_{L=1}^{\infty} \frac{1}{L!} \left(\frac{ik\mu}{aH} \right) T_{//}^{L}(\mathbf{k}) \quad \text{where} \quad T_{//}^{L}(\mathbf{x}) = \left[1 + \delta(\mathbf{x}) \right] u^{L}(\mathbf{x})$$

Redshift-space momentum field can also be described by a sum over real-space velocity moments.

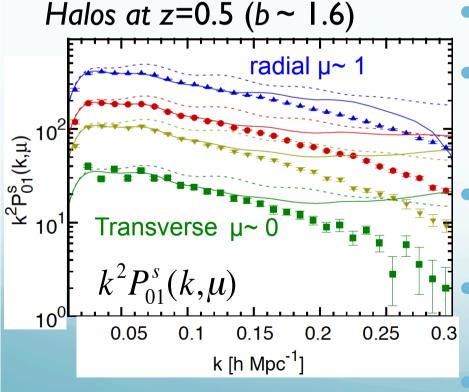
$$p_{\parallel}^{s}(\mathbf{x}) = \left[1 + \delta^{s}(\mathbf{x})\right] v_{\parallel}^{s}(\mathbf{x})$$



Expression for density-momentum power

$$\langle \delta^{s}(\mathbf{k}) p_{//}^{s^{*}}(\mathbf{k}') \rangle = (2\pi)^{3} P_{01}^{s}(\mathbf{k}) \delta^{D}(\mathbf{k} - \mathbf{k}')$$

$$-i\frac{k\mu}{\mathcal{H}}P_{01}^{s}(\mathbf{k}) = -i\frac{k\mu}{\mathcal{H}}P_{01}^{r}(\mathbf{k}) + \left(\frac{k\mu}{\mathcal{H}}\right)^{2}P_{11}^{r}(\mathbf{k}) - \left(\frac{k\mu}{\mathcal{H}}\right)^{2}P_{02}^{r}(\mathbf{k}) + \frac{i}{2}\left(\frac{k\mu}{\mathcal{H}}\right)^{3}P_{03}^{r}(\mathbf{k})$$
$$-2i\left(\frac{k\mu}{\mathcal{H}}\right)^{3}P_{12}^{r}(\mathbf{k}) + \frac{1}{6}\left(\frac{k\mu}{\mathcal{H}}\right)^{4}P_{04}^{r}(\mathbf{k}) - \left(\frac{k\mu}{\mathcal{H}}\right)^{4}P_{13}^{r}(\mathbf{k}) + \frac{1}{2}\left(\frac{k\mu}{\mathcal{H}}\right)^{4}P_{22}^{r}(\mathbf{k})$$

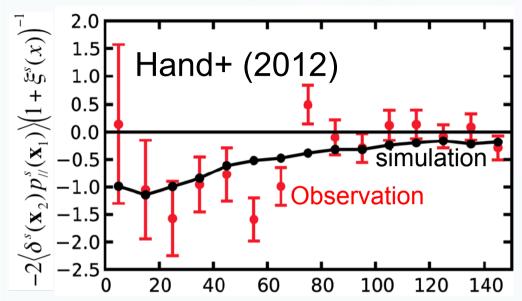


- Data points are simulation results
- Dotted lines are linear theory (Kaiser)

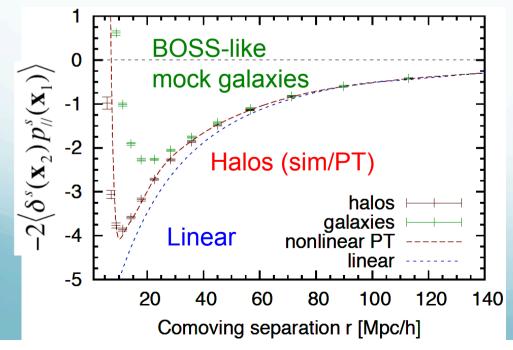
$$P_{01}^{s}(k,\mu) = (if\mu/k)(b + f\mu^{2})P_{m}(k)$$

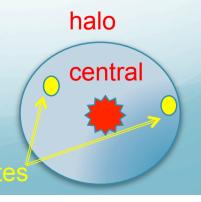
- Solid lines are 1-loop PT + non-local bias + DM spectrum from sim
 - Suppression starting at large scales is nonlinear velocity dispersion (RSD)
 - No free parameter for RSD

Redshift-space pairwise infall momentum



- $-2\langle \delta^s(\mathbf{x}_2) p_{\parallel}^s(\mathbf{x}_1) \rangle$: Fourier tr of P_{01}
- Transition of sign seen in observed data on small scales is caused by nonlinear RSD.
- PT well explains the halo results.
- Finger-of-god for satellites needs to be taken into account.



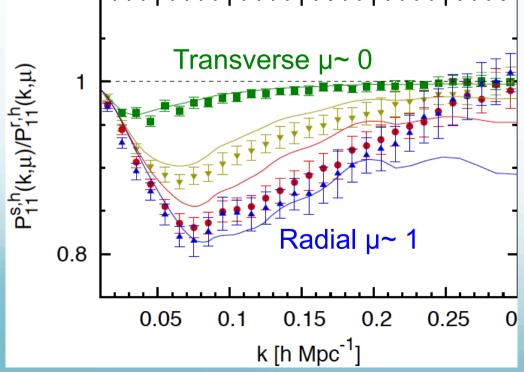


Momentum-momentum power spectrum

$$\langle p_{\parallel}^{s}(\mathbf{k})p_{\parallel}^{s*}(\mathbf{k}')\rangle = (2\pi)^{3}\underline{P_{11}^{s}(\mathbf{k})}\delta^{D}(\mathbf{k}-\mathbf{k}')$$

$$\left(\frac{k\mu}{\mathcal{H}}\right)^{2} P_{11}^{s}(\mathbf{k}) = \left(\frac{k\mu}{\mathcal{H}}\right)^{2} P_{11}^{r}(\mathbf{k}) - 2i\left(\frac{k\mu}{\mathcal{H}}\right)^{3} P_{12}^{r}(\mathbf{k}) - \left(\frac{k\mu}{\mathcal{H}}\right)^{4} P_{13}^{r}(\mathbf{k}) + \left(\frac{k\mu}{\mathcal{H}}\right)^{4} P_{22}^{r}(\mathbf{k})$$

Ratio of redshift space and real space for halos P_{11}^{s}/P_{11}^{r}



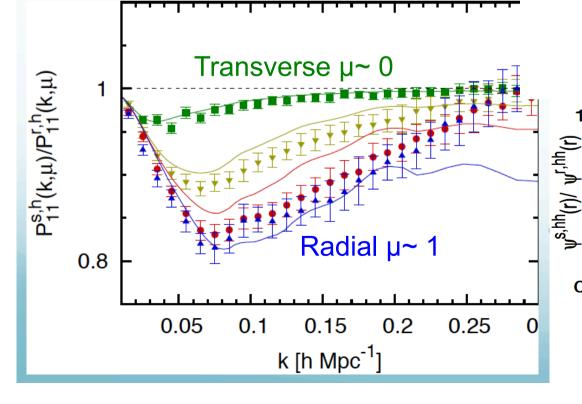
- Unity in linear theory
 - Deviation from unity is entirely due to nonlienar RSD (~20%)
- Affected by shot noise

Momentum-momentum power spectrum

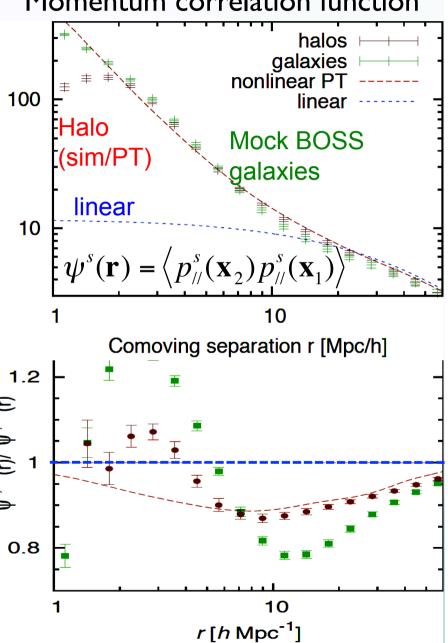
$$\langle p''_{\parallel}(\mathbf{k}) p''_{\parallel}(\mathbf{k}') \rangle = (2\pi)^3 \underline{P_{11}^s(\mathbf{k})} \delta^D(\mathbf{k} - \mathbf{k}')$$

$$\left(\frac{k\mu}{\mathcal{H}}\right)^{2} P_{11}^{s}(\mathbf{k}) = \left(\frac{k\mu}{\mathcal{H}}\right)^{2} P_{11}^{r}(\mathbf{k}) - 2i\left(\frac{k\mu}{\mathcal{H}}\right)^{2} - \left(\frac{k\mu}{\mathcal{H}}\right)^{4} P_{13}^{r}(\mathbf{k}) + \left(\frac{k\mu}{\mathcal{H}}\right)^{4} P_{0}^{s}$$

Ratio of redshift space and real space P_1



Momentum correlation function



Conclusions

- Observables from kSZ survey is momentum field (not velocity field) in redshift-space.
- Expressions for velocity statistics in redshift space were derived.
- Power spectra were computed based on I-loop nonlinear PT and compared to N-body results for halos/dark matter.
- Linear theory does not work even at very large scales.
- See Okumura, Seljak, Vlah & Desjacques, (arXiv:1312.4214) for the detail

To improve precision

- Better understanding for shot noise is required for momentum autopower spectrum P_{11}
- We have a good model for halos, but not yet for galaxies (I halo term, Finger-of-God, halo models, etc).