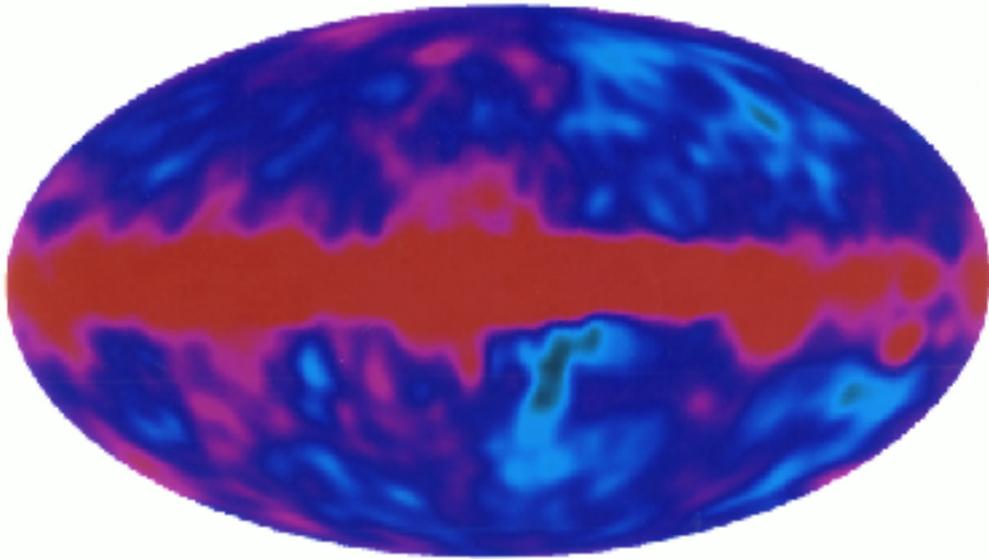


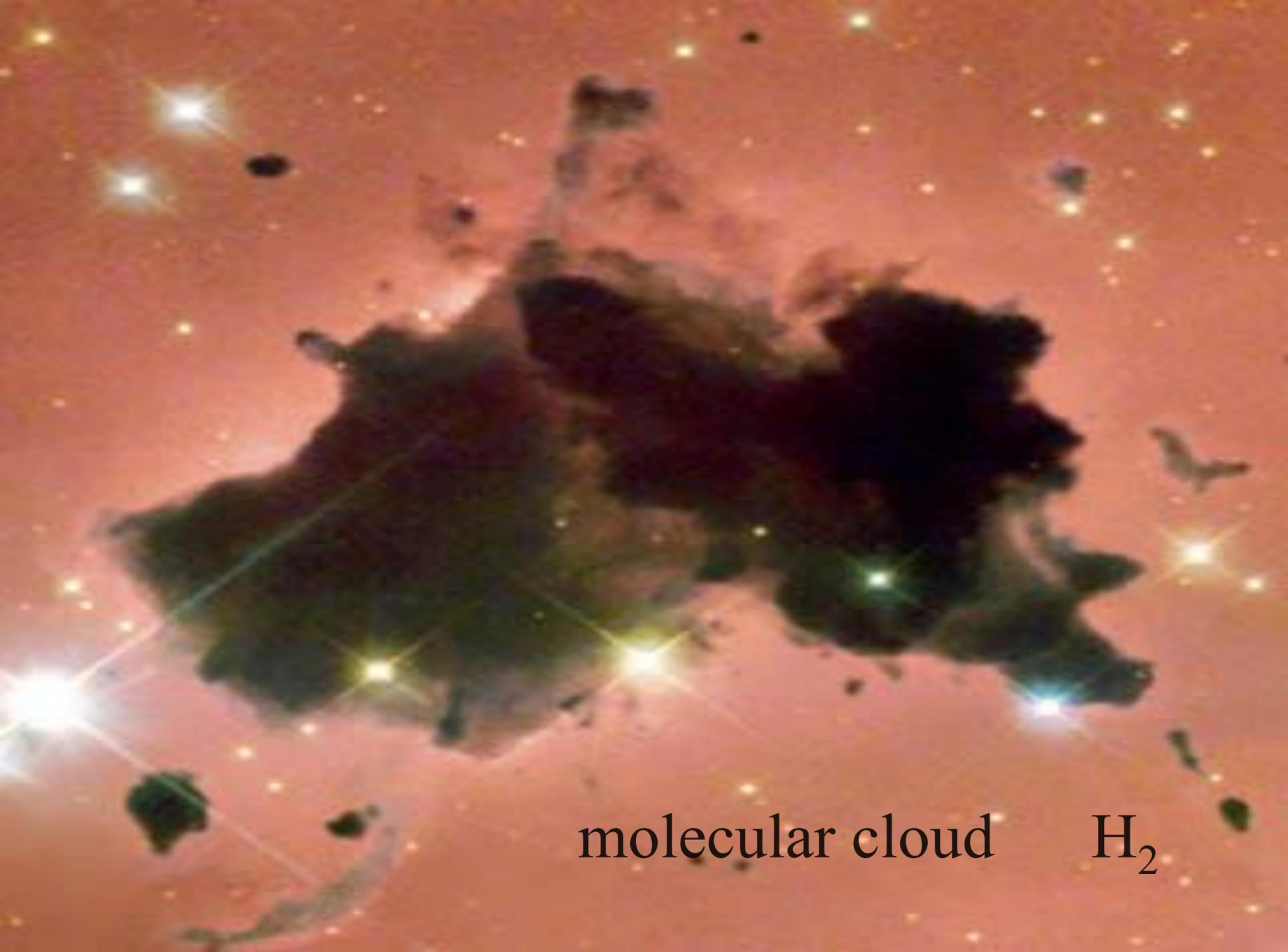
Struggling with Entropy

Charley Lineweaver
Research School of Astronomy and Astrophysics
Research School of Earth Sciences
Australian National University



CMB and Aliens





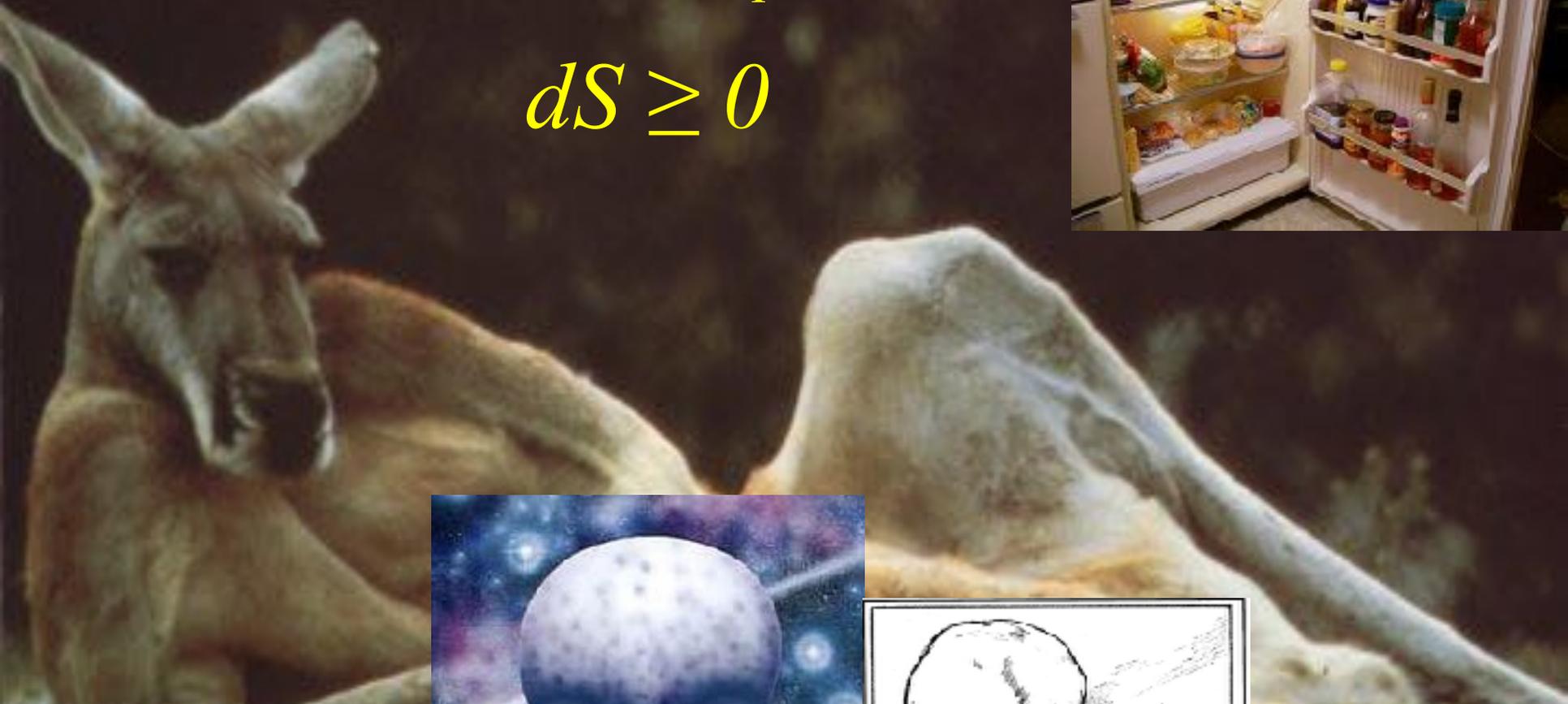
molecular cloud H_2



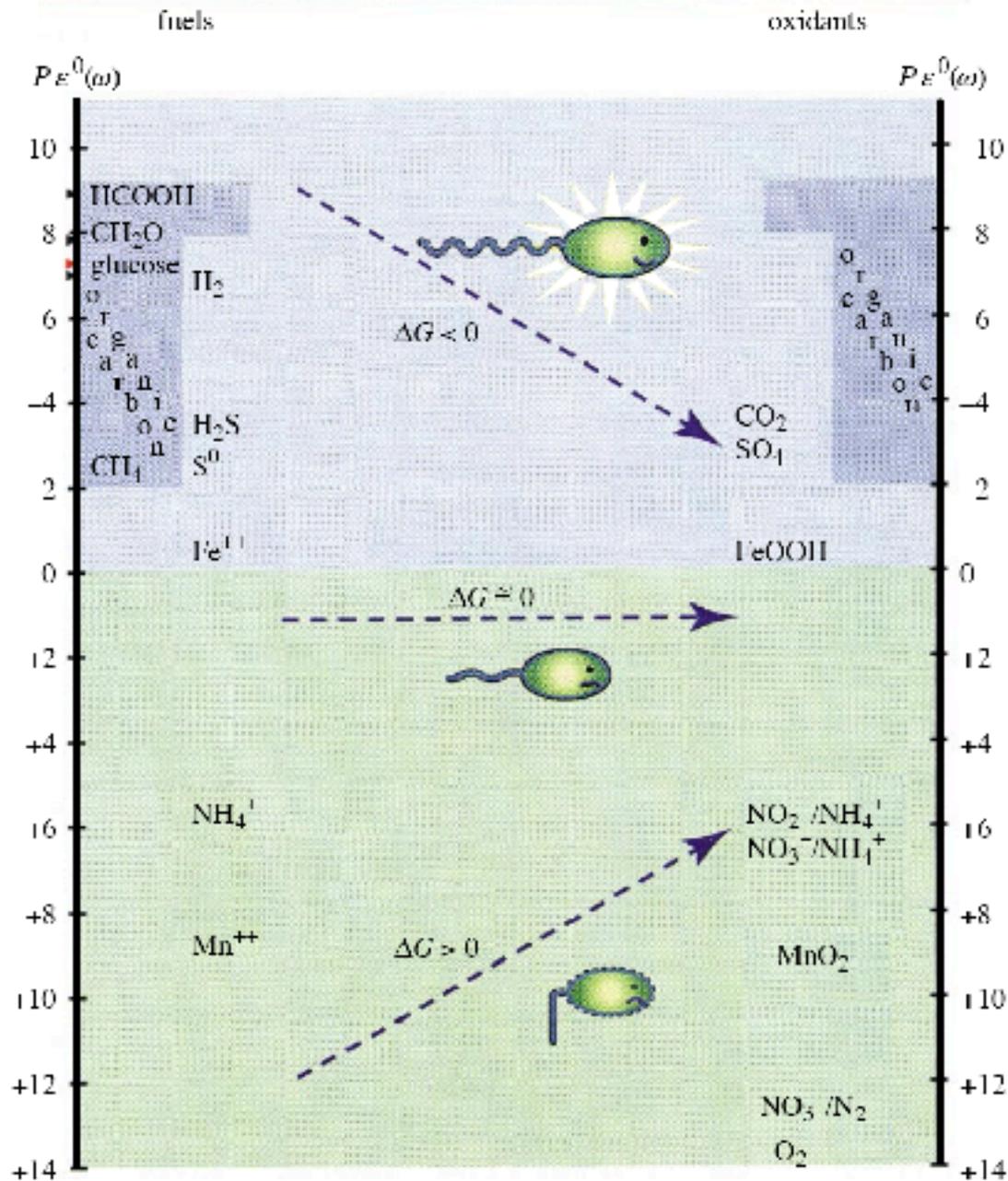
life form CHON

Life forms, refrigerators
and all other irreversible processes

$$dS \geq 0$$



thermodynamics
it's not just a good idea... it's the law!



Prochlorococcus
100 octillion (10²⁹) individuals



Nealson and Conrad 1999

**Is the complexity of
the universe increasing?**



**COMPLEXITY
AND THE
ARROW OF TIME**

Edited by

Charles H. Lineweaver,
Paul C. W. Davies and
Michael Ruse



**Is the complexity of
the universe increasing?**

**Many definitions of complexity
but all of them require free energy.
All require an entropy gap ΔS
between the entropy of the universe
and the maximum entropy of the universe.**

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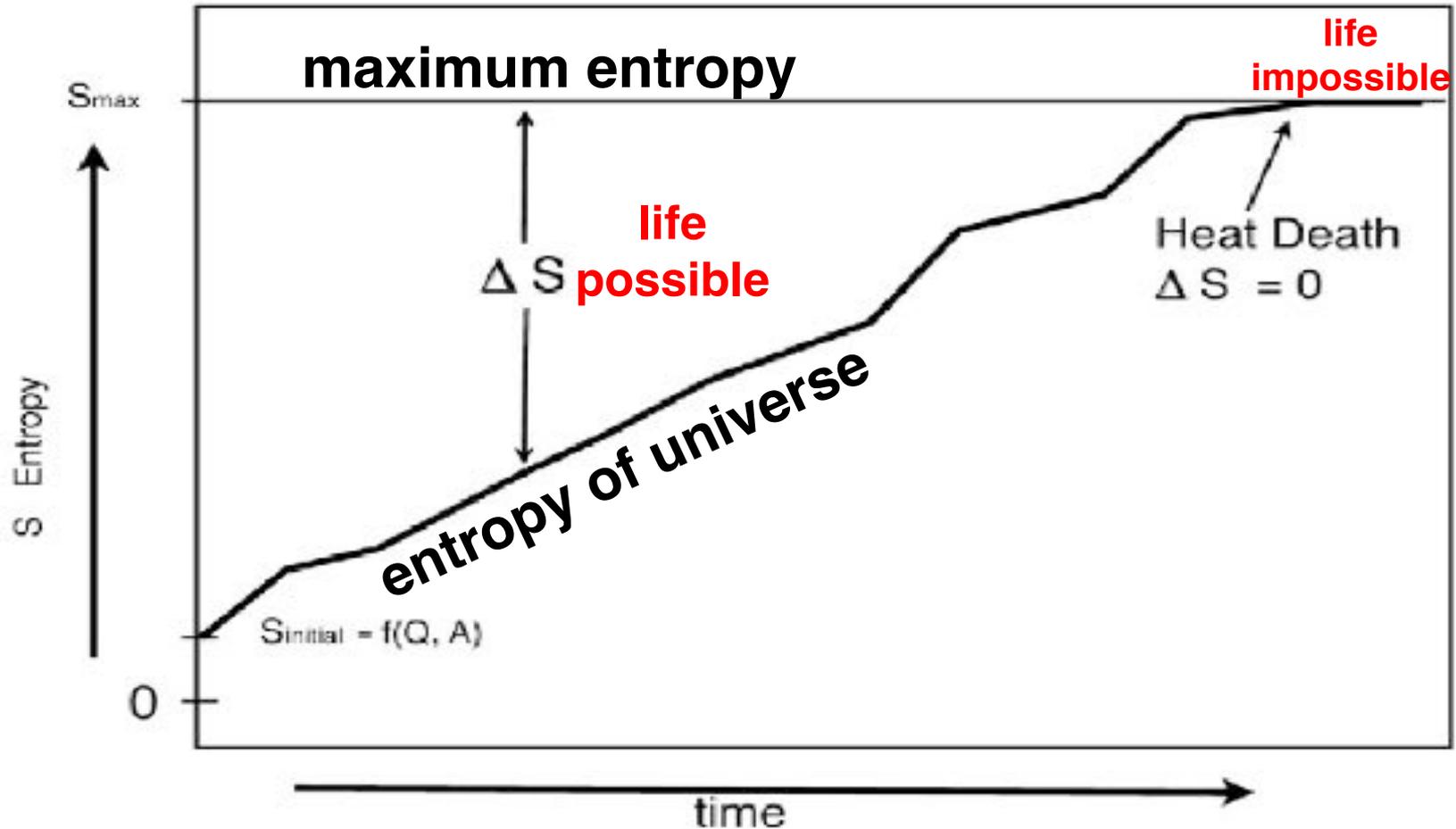
**How can complexity keep
increasing if we are headed
for a heat death?**

**COMPLEXITY
AND THE
ARROW OF TIME**

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Second law of cosmobiology



Life and all other dissipative structures in the universe depend on the non-equilibrium of the universe and the availability of free energy. Within the standard Λ CDM cosmology, we quantify this non-equilibrium as the difference between the maximum entropy the universe will ever have (within the comoving volume of the current particle horizon), and the monotonically increasing entropy of the universe within the same comoving volume. When $S_{max} - S_{uni}(t) \approx 0$ the universe will reach an effective heat death from which no more free energy can be extracted.

A LARGER ESTIMATE OF THE ENTROPY OF THE UNIVERSE

CHAS A. EGAN

Research School of Astronomy and Astrophysics, Australian National University, Canberra, Australia ¹

CHARLES H. LINEWEAVER

Planetary Science Institute, Research School of Astronomy and Astrophysics and Research School of Earth Sciences, Australian National University, Canberra, Australia

Received 2009, September 22; accepted 2010, January 11.

ABSTRACT

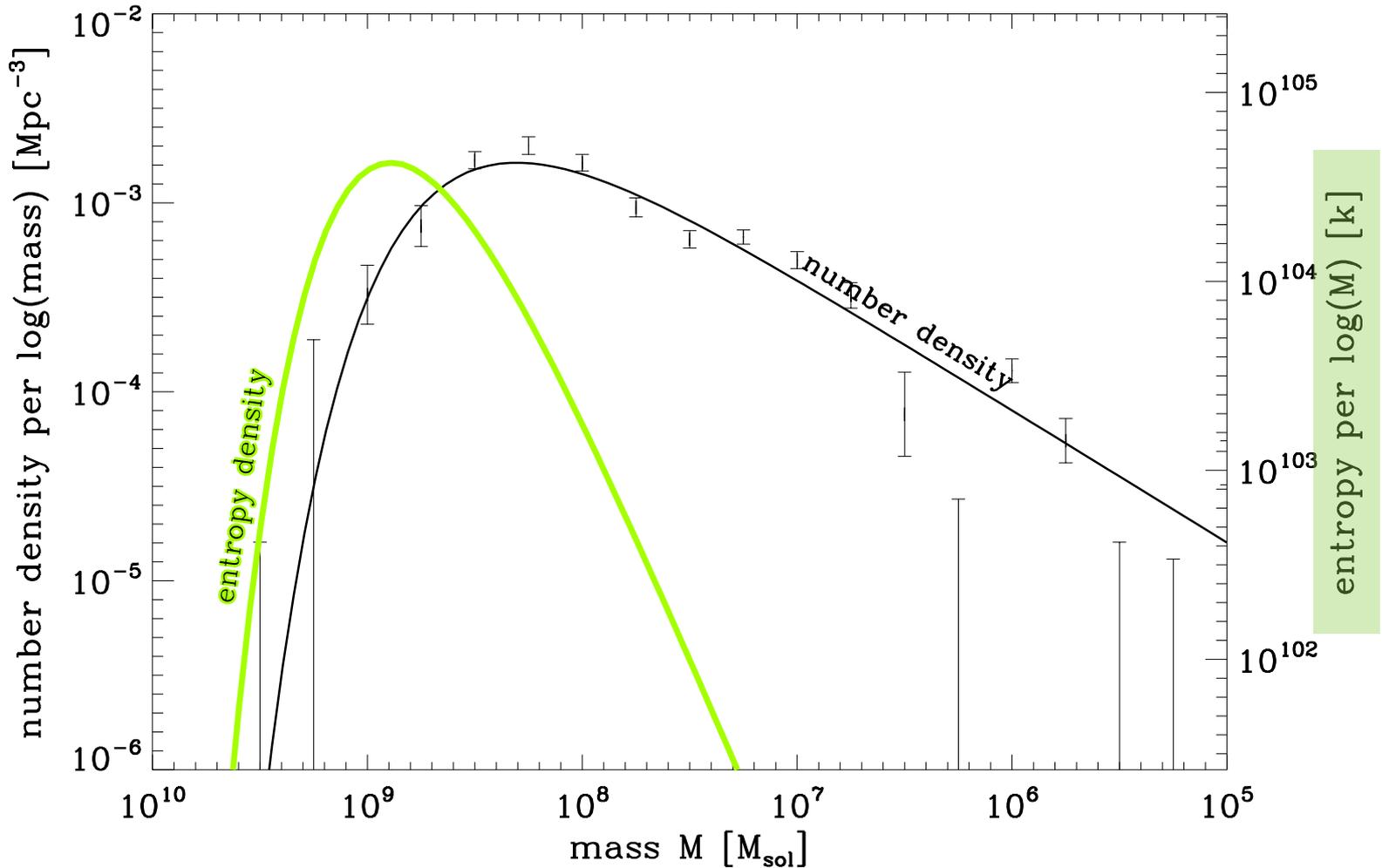
Using recent measurements of the supermassive black hole (SMBH) mass function, we find that SMBHs are the largest contributor to the entropy of the observable universe, contributing at least an order of magnitude more entropy than previously estimated. The total entropy of the observable universe is correspondingly higher, and is $S_{\text{obs}} = 3.1_{-1.7}^{+3.0} \times 10^{104} k$. We calculate the entropy of the current cosmic event horizon to be $S_{\text{CEH}} = 2.6 \pm 0.3 \times 10^{122} k$, dwarfing the entropy of its interior, $S_{\text{CEH int}} = 1.2_{-0.7}^{+1.1} \times 10^{103} k$. We make the first tentative estimate of the entropy of weakly interacting massive particle dark matter within the observable universe, $S_{\text{dm}} = 10^{88} \pm 1 k$. We highlight several caveats pertaining to these estimates and make recommendations for future work.

*The initial low gravitational entropy of the universe as the origin of design in nature.
Lineweaver and Egan 2011 (Chapter in Origin of Design, Edt Seckbach et al, Springer*

Component	Entropy S [k] (previous work)	Entropy S [k]	Entropy Density s [$k m^{-3}$]
SMBHs S_{SMBH}	10^{101} [26], 10^{102} [3]	$2.9_{-1.6}^{+2.8} \times 10^{104}$	$8.4_{-4.7}^{+8.2} \times 10^{23}$
Stellar BHs S_{SBH}	10^{97} [3], 10^{98} [1]	$5.6 \times 10^{97-0.6}$	$1.6 \times 10^{17-0.6}$
CMB S_γ	10^{88} [1, 3, 26], 10^{89} [34]	$5.104 \pm 0.011 \times 10^{89}$	$1.478 \pm 0.003 \times 10^9$
Relic Neutrinos S_ν	10^{88} [3], 10^{89} [34]	$4.87 \pm 0.01 \times 10^{89}$	$1.41 \pm 0.01 \times 10^9$
Dark Matter S_{DM}	—	$1.0 \times 10^{88 \pm 1}$	$2.9 \times 10^{7 \pm 1}$
Relic Gravitons S_g	10^{86} [3]	$2.8 \times 10^{86 \pm 1.3}$	$1.2 \times 10^{6 \pm 1.3}$
Non-Relic Photons $S_{\gamma_{nr}}$	10^{86} [1, 3, 35]	—	—
ISM & IGM S_{gas}	—	$4 \times 10^{81 \pm 0.5}$	$1.2 \times 10^{1 \pm 0.5}$
Stars S_*	10^{79} [3]	$1.2 \times 10^{80 \pm 0.4}$	$3.5 \times 10^{-1 \pm 0.4}$
Total S_{uni}	10^{101} [26], 10^{102} [3]	$2.9_{-1.6}^{+2.8} \times 10^{104}$	$8.4_{-4.7}^{+8.2} \times 10^{23}$

- [1] Frautschi 1982
[3] Frampton et al 2008
[26] Penrose 2004
[34] Kolb & Turner 1981
[35] Bousso et al 2007

Black hole mass function



$$S_{BH} \propto k M^2$$

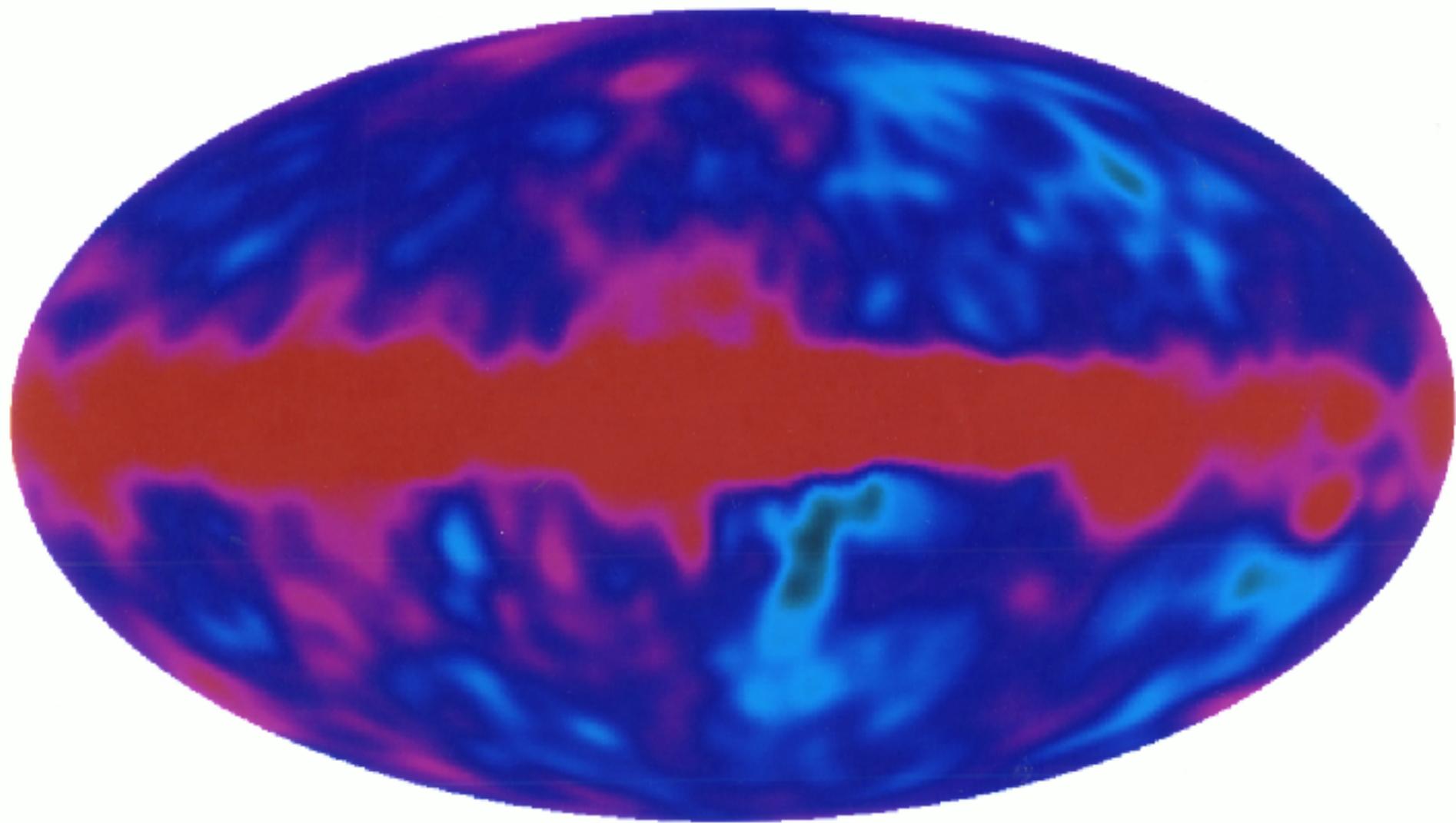


me

George
Smoot

May 2, 1992

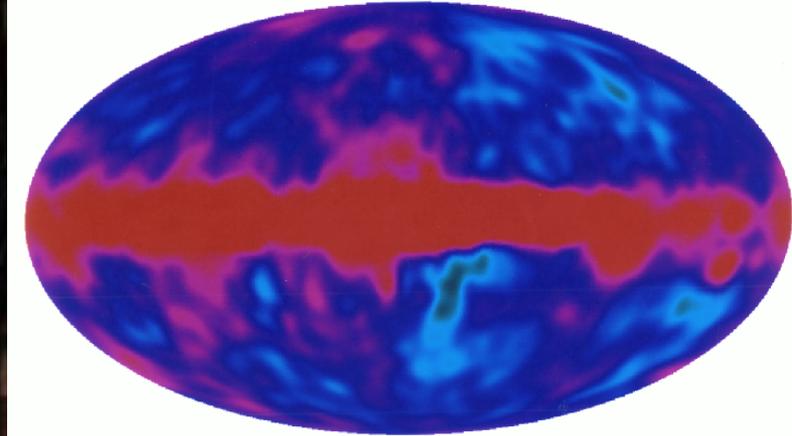
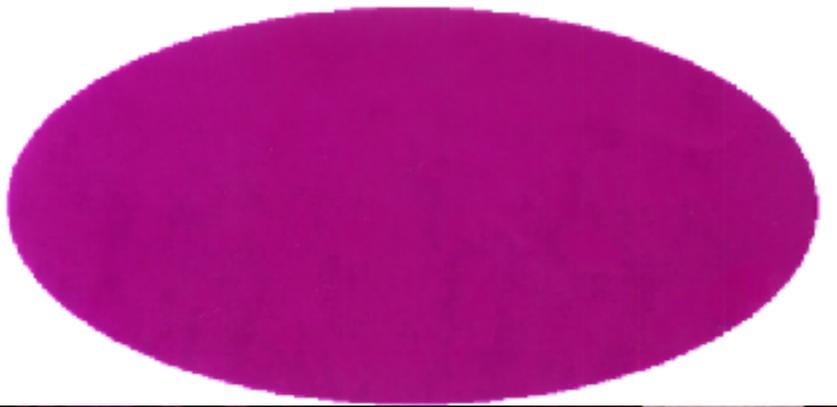
188



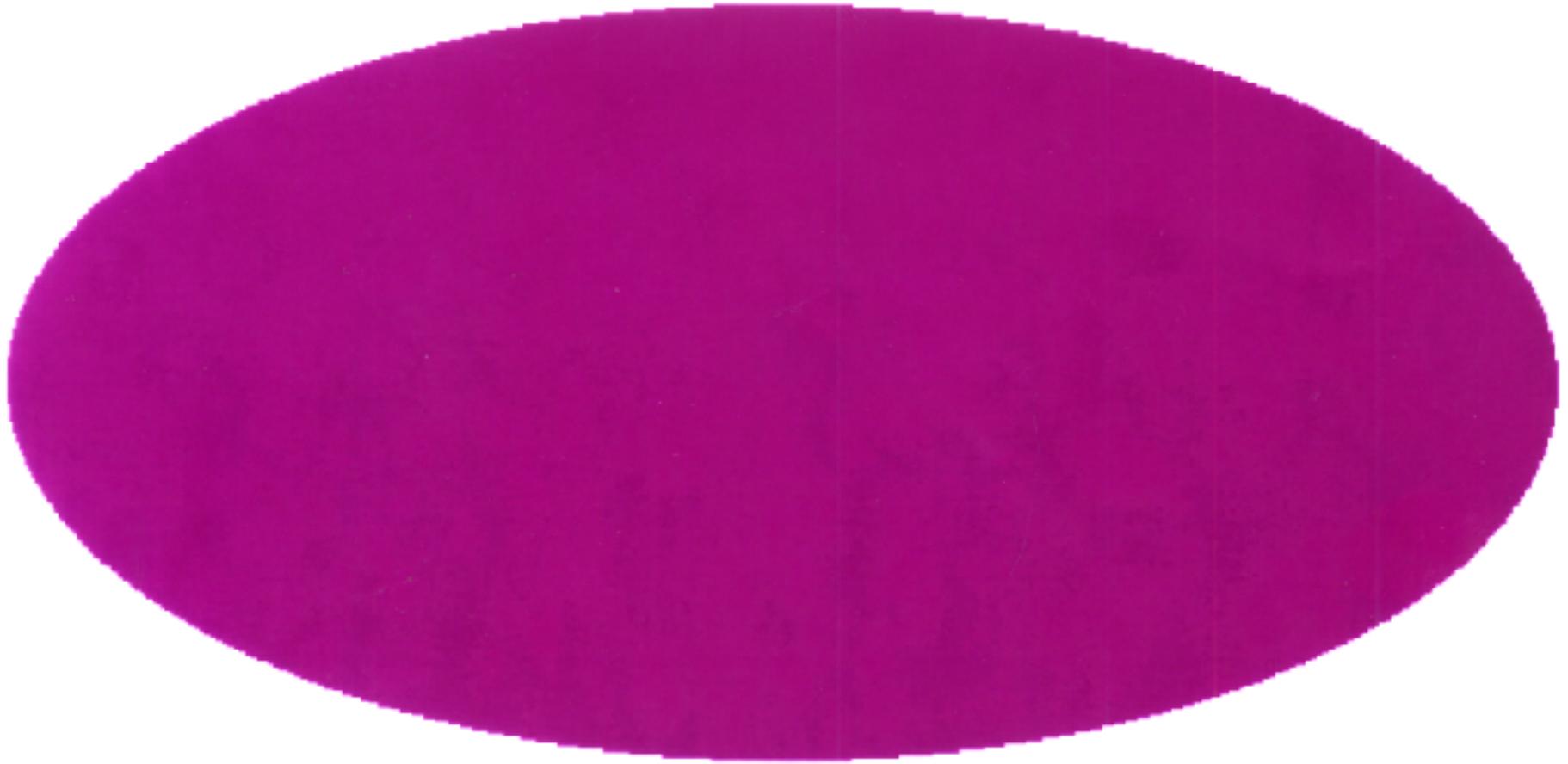
2006 Physics Nobel Prize

John Mather

George Smoot



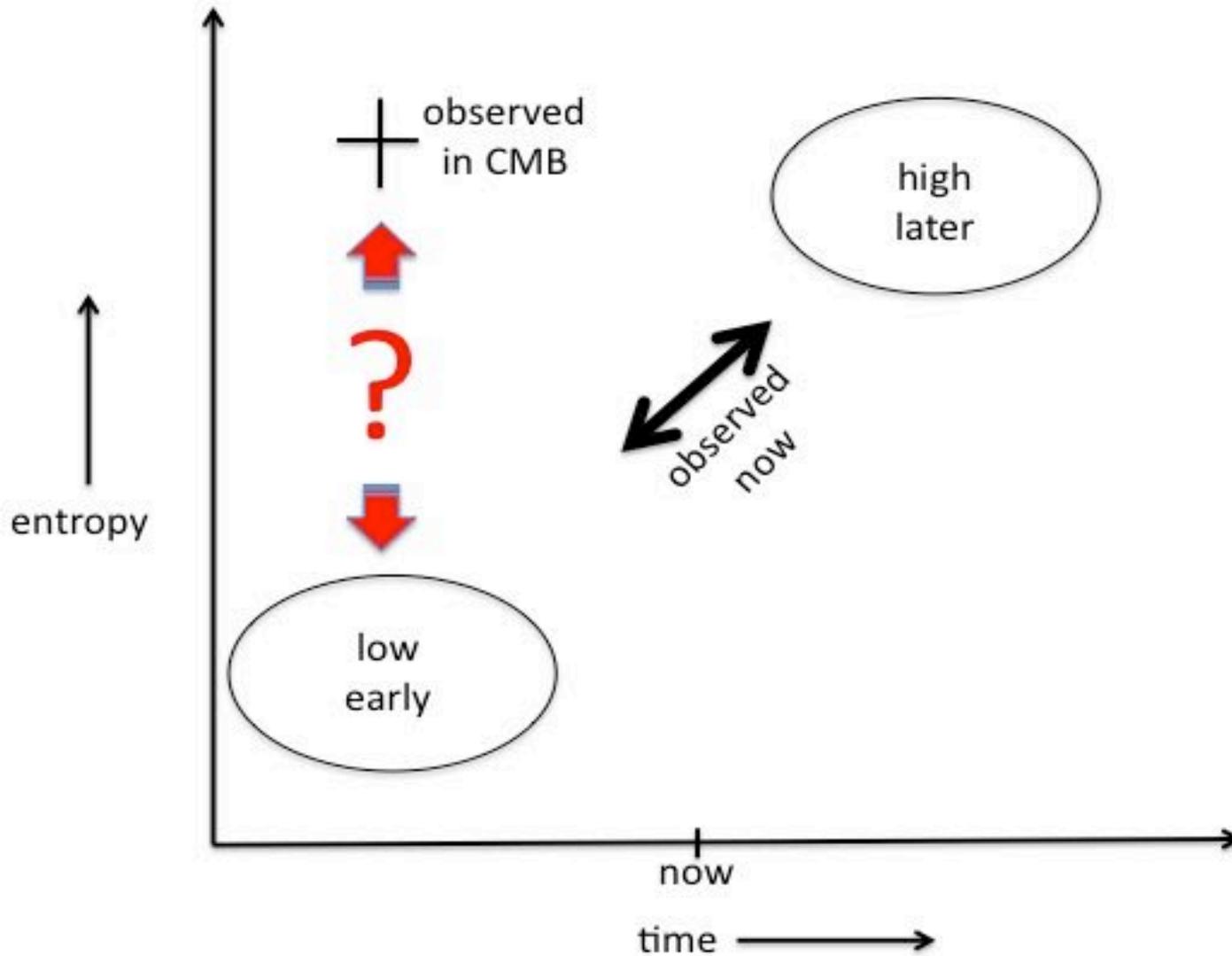
The universe 380,000 years after the big bang.



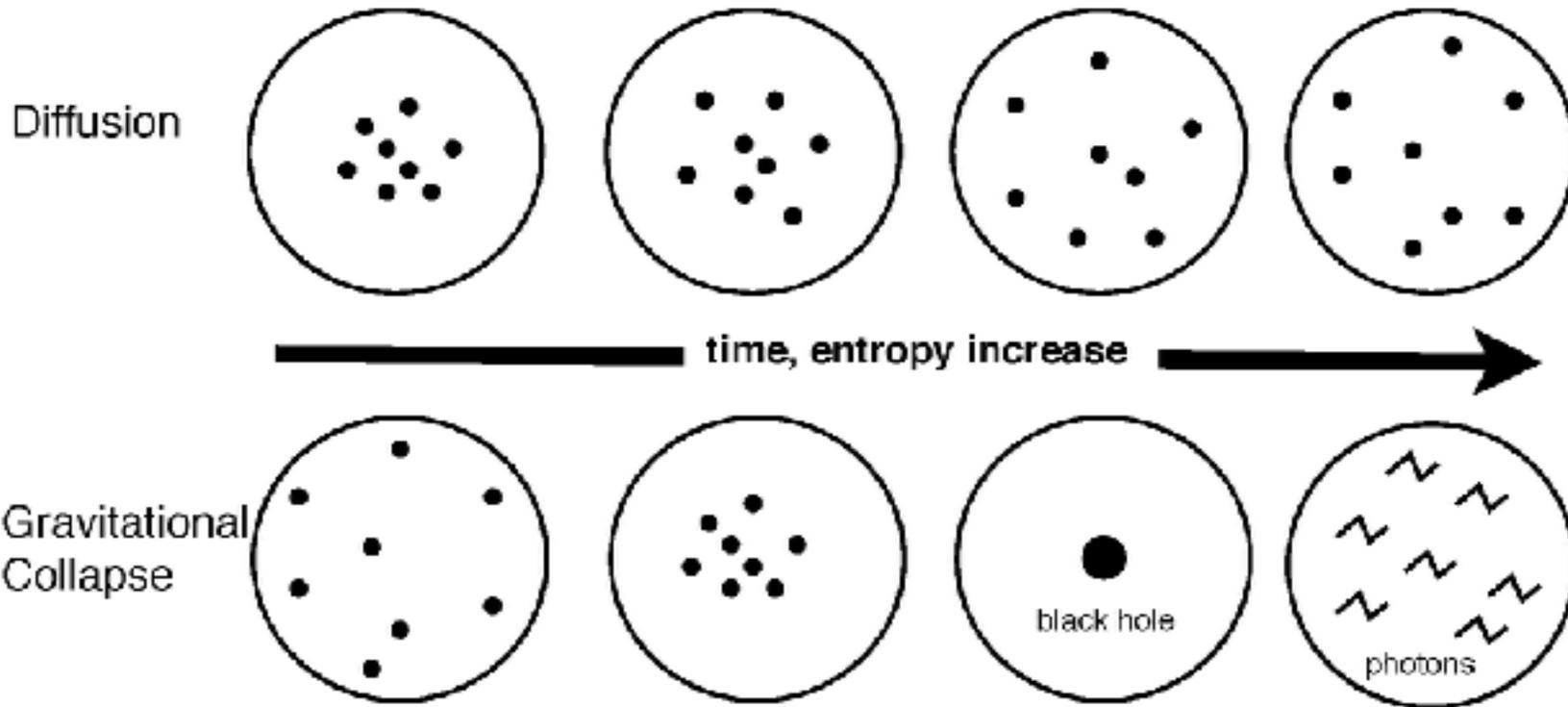
Full Sky Map of the Cosmic Microwave Background

COBE FIRAS and DMR instruments

The Initial Entropy Problem

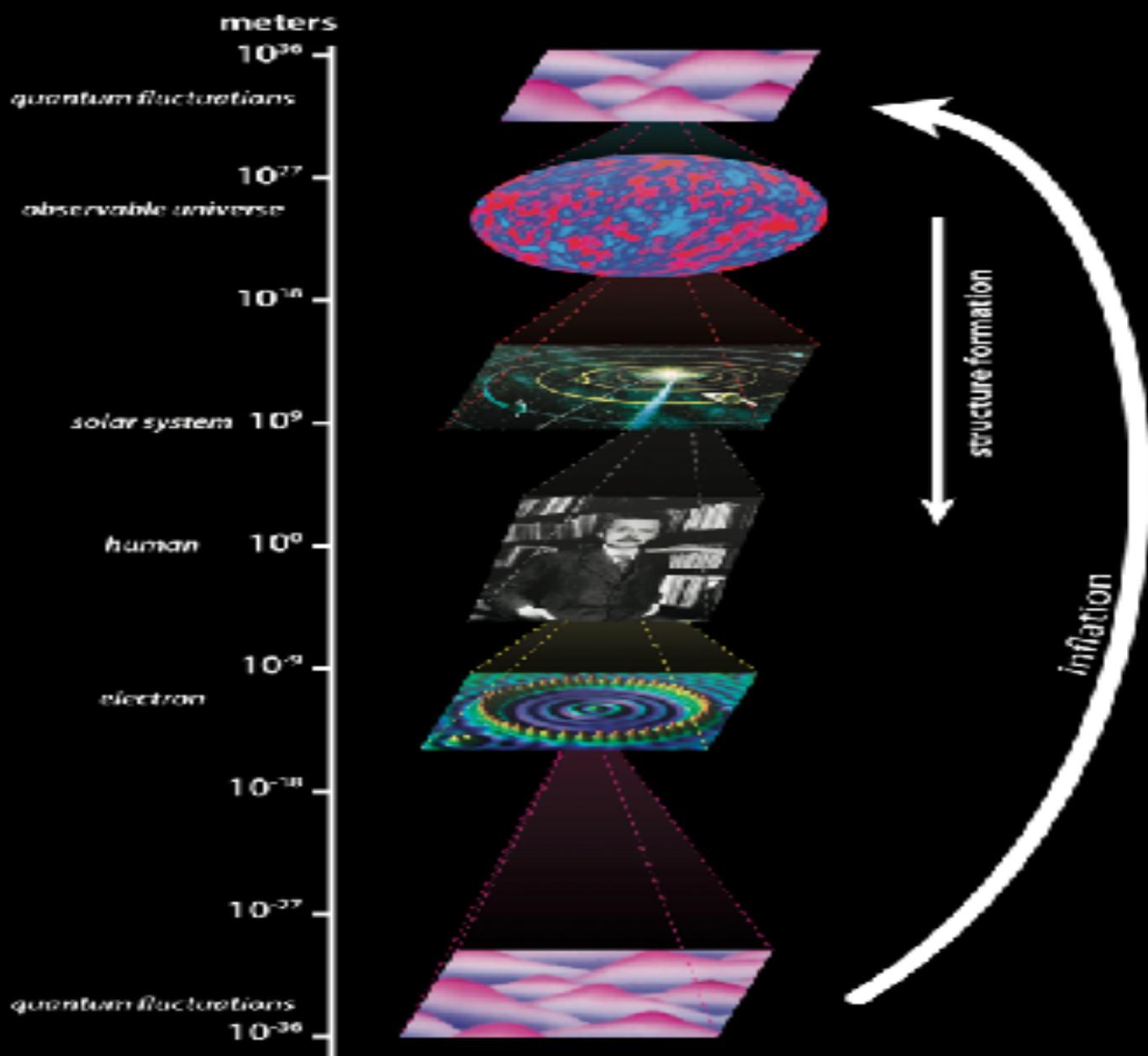


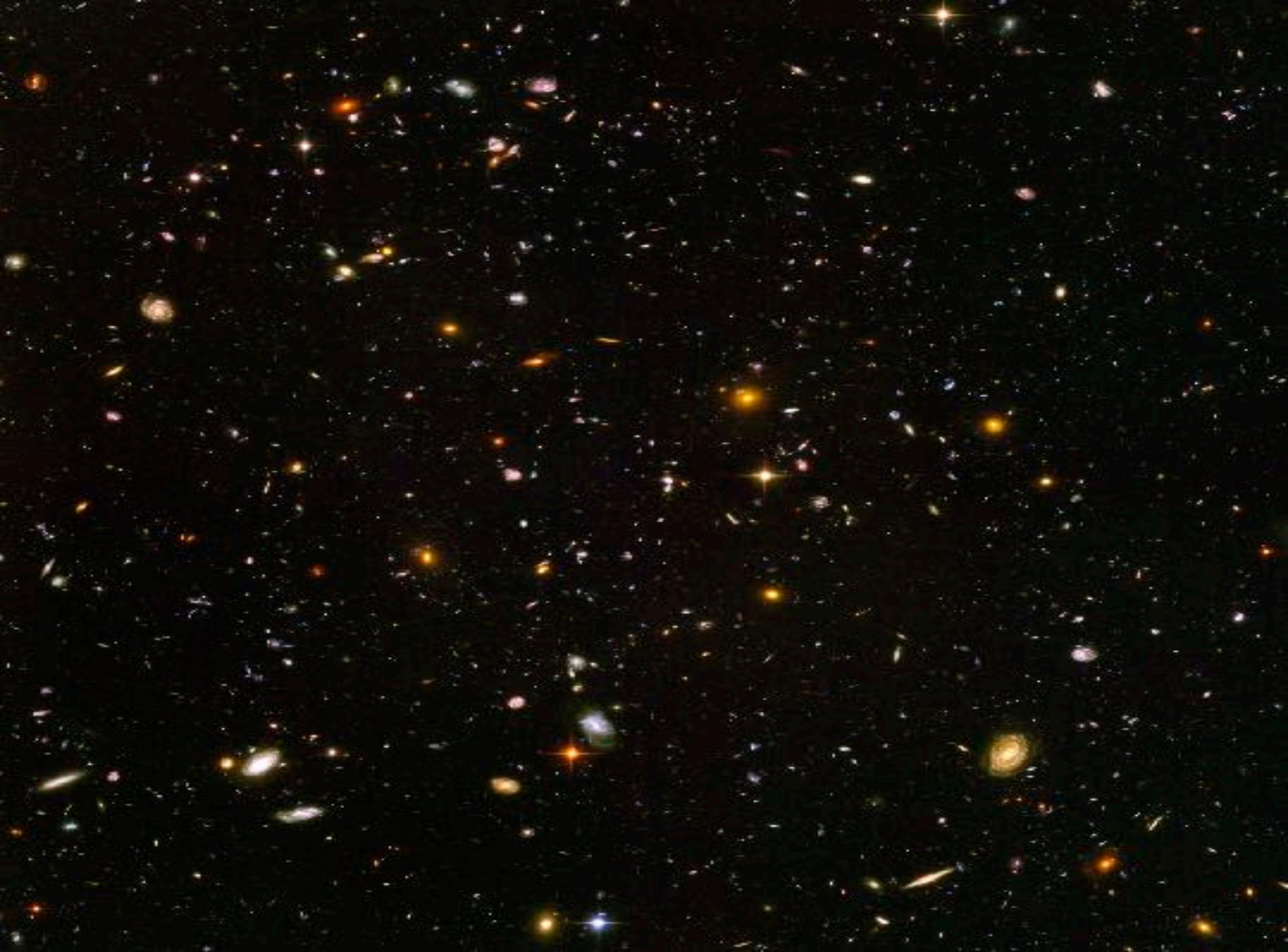
kinetic energy dominated system: heat capacity > 0



gravity dominated system: heat capacity < 0

Origin of Structure





Beware “temporal provincialism*”:

The tendency to slip in assumptions about $\dot{S} > 0$
(and thus tunings of initial conditions)
without even realizing it

Related issues:

- Arrival Terminals

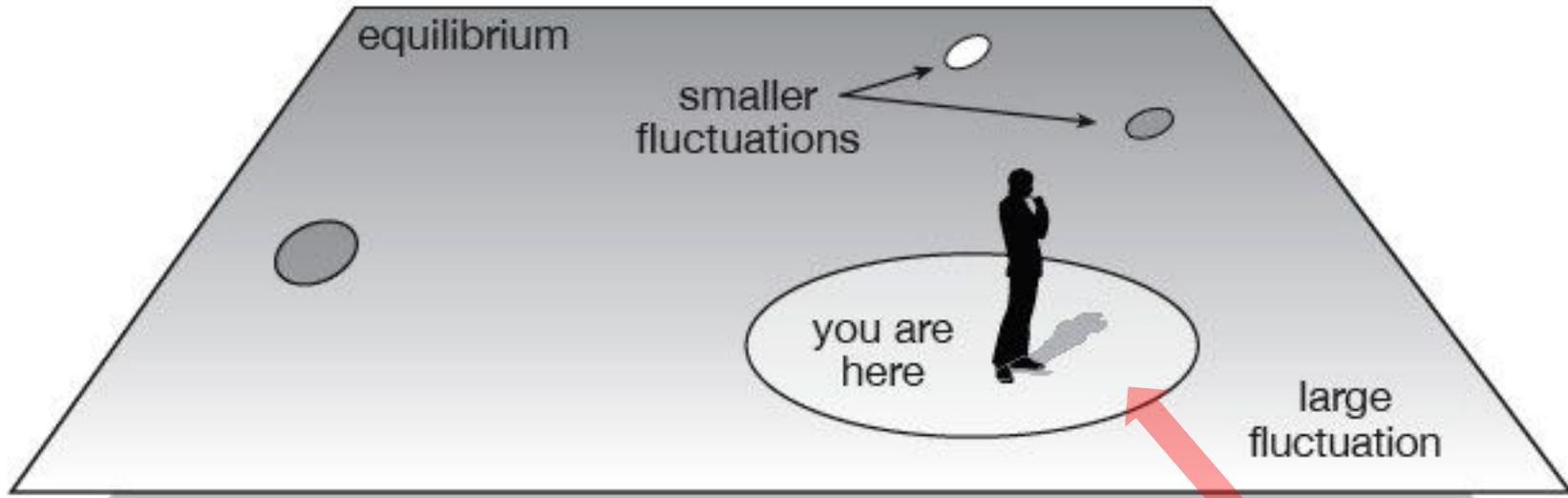
* L. Susskind

**"Temporal Provincialism" ?
you can't argue with a cup of coffee**

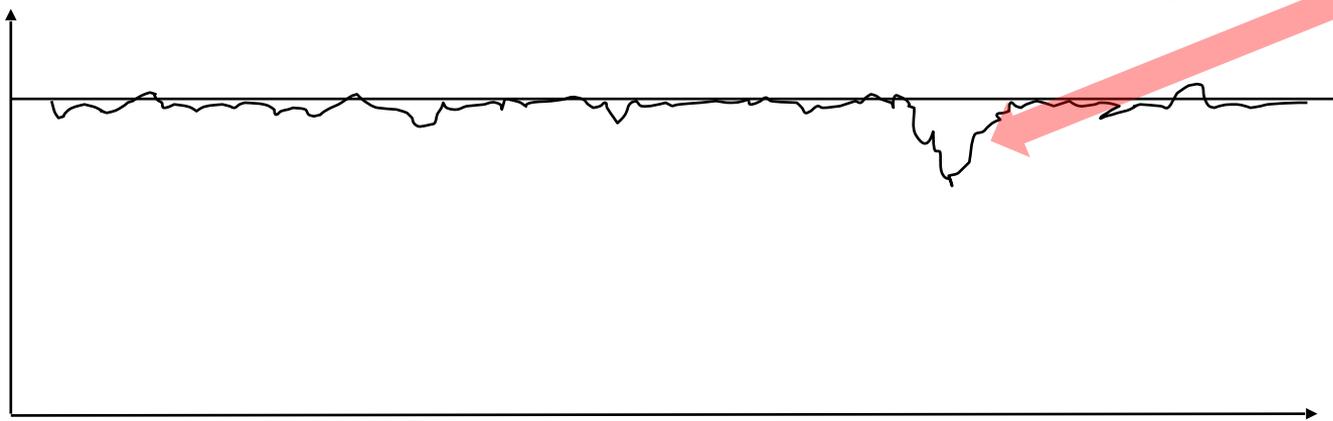


$$dS \geq 0$$

**the low entropy of our universe interpreted as
a fluctuation away from an average equilibrium universe**



**minimum fluctuation from equilibrium
(spatially and temporally)
compatible with our existence**



time

$V(t_0) = \text{red line}$

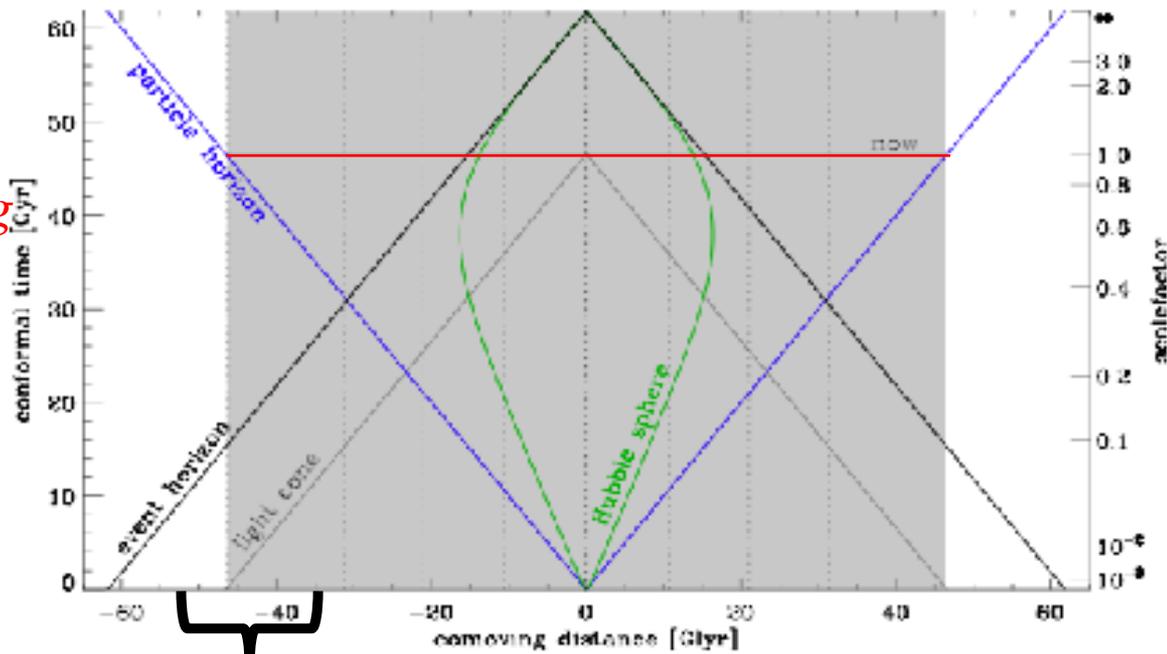
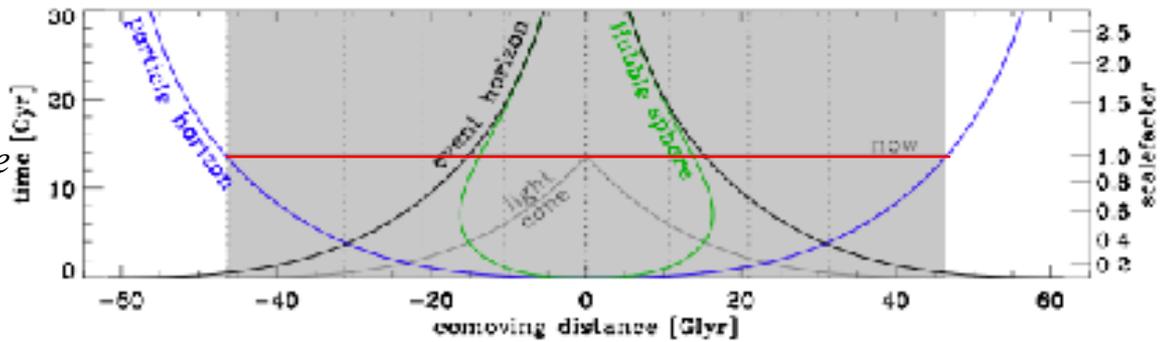
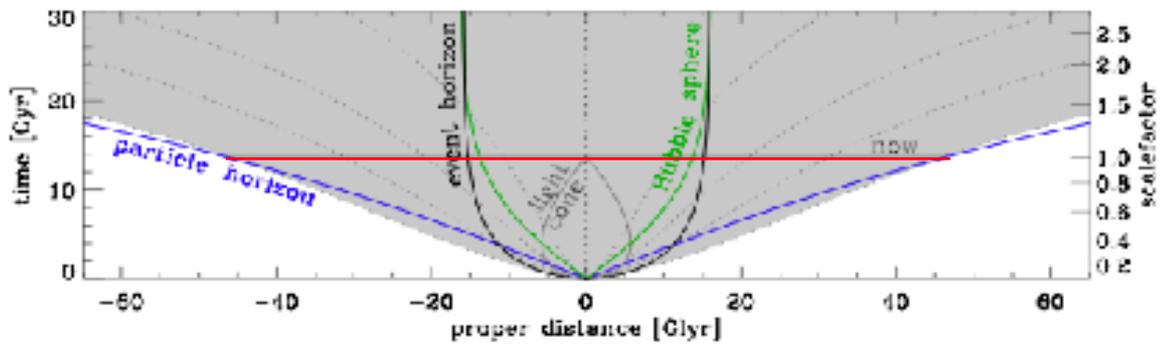
$V(t) = \text{grey area}$

= comoving volume of the current particle horizon

Λ CDM

NB: V is larger than the observable universe in the past and smaller than the observable universe in the future.

as time goes by,
we see more of the universe
which seems to be just like the
universe we have been observing



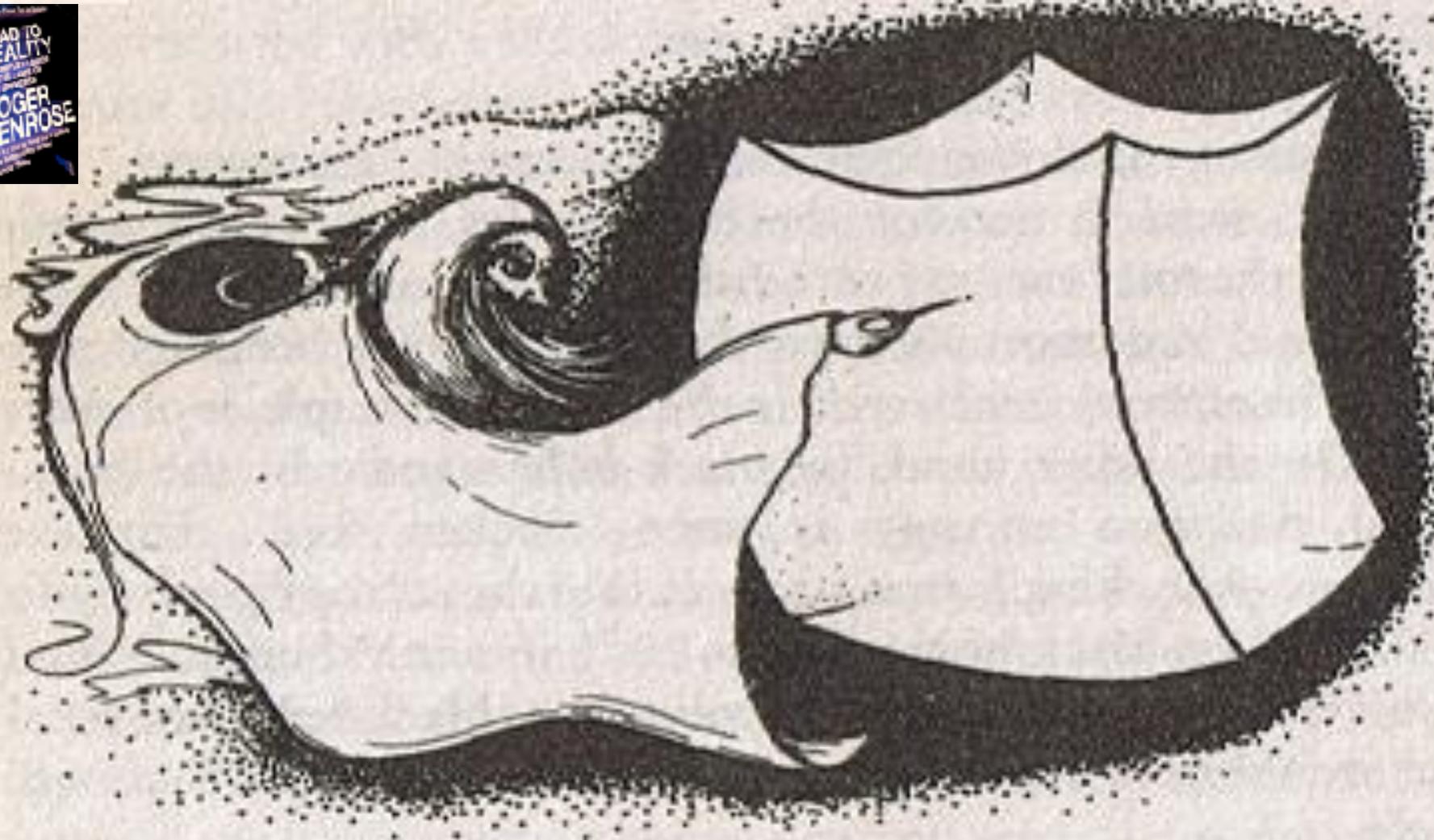


Fig. 7.19. In order to produce a universe resembling the one in which we live, the Creator would have to aim for an absurdly tiny volume of the phase space of possible universes – about $1/10^{10^{123}}$ of the entire volume, for the situation under consideration. (The pin, and the spot aimed for, are not drawn to scale!)

If you have to invoke fine-tuning to pick out a low entropy initial condition from a gigantic equilibrium-dominated phase space, why not dispense with the imagined phase space and just start with the most natural low entropy initial condition you can think of.

(cut out the middle man)

S_{max}

maximum entropy

S Entropy

0

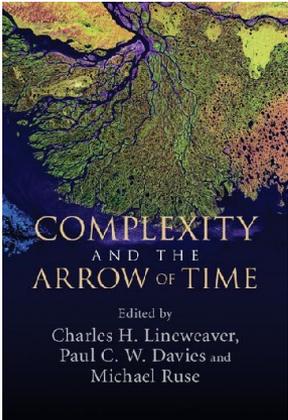
$S_{initial} = f(Q, A)$

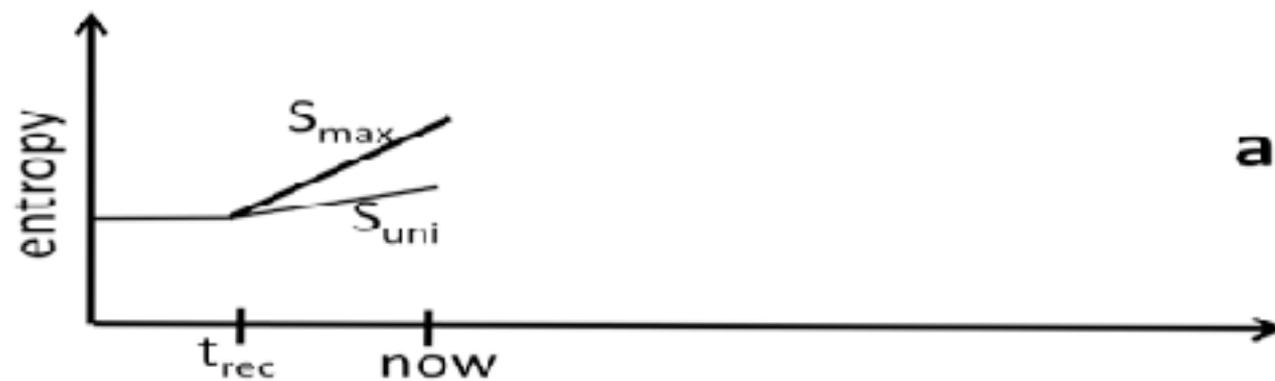
entropy of universe

ΔS

Heat Death
 $\Delta S = 0$

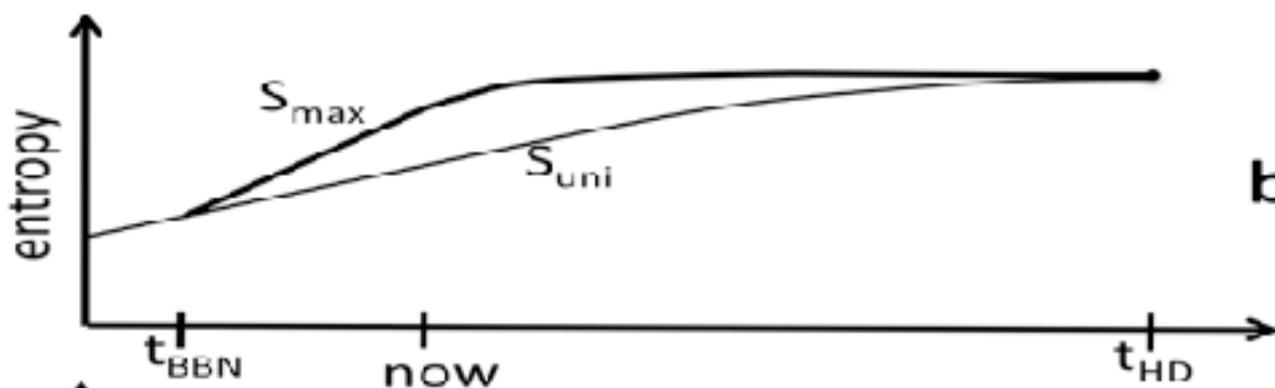
time





a

Layzer 1975
 Frautschi 1988
 Barrow 1994
 Chaisson 2001



b

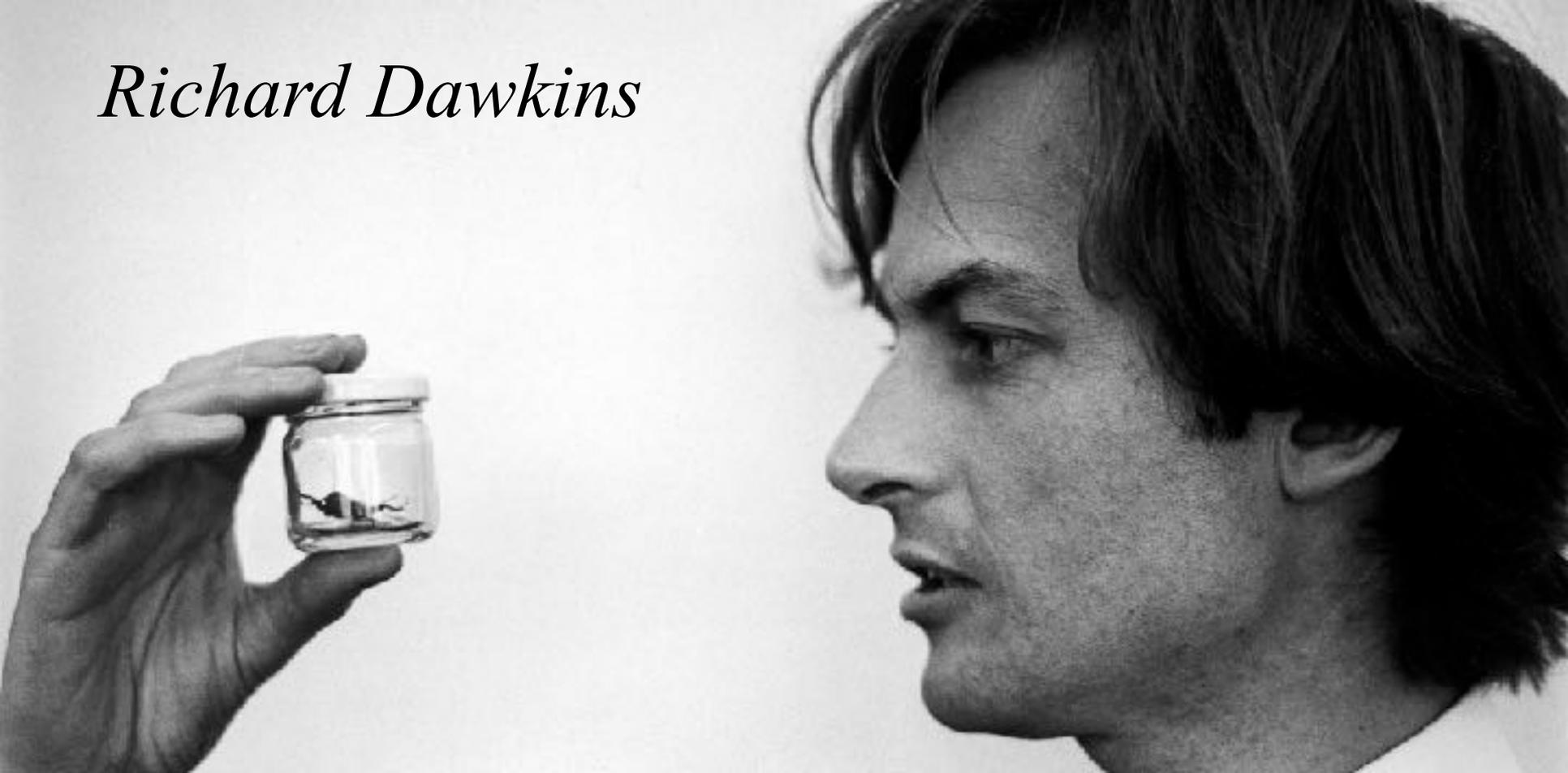
Davies 1994
 Barrow 2011



c

Penrose 2004
 Lineweaver & Egan 2008

Richard Dawkins



*We, the living, should feel lucky.
Just think of all the people who haven't been born.*

$$P = N_{\text{living}} / N_{\text{non-living}}$$

$$P \sim 0$$

**To determine probability, or luck, or
"tuning"**

you need to know:

**1) what has been selected
and**

2) what it has been selected from

**ie you need to know what is and
what could have been.**

Contrafactuals are problematic

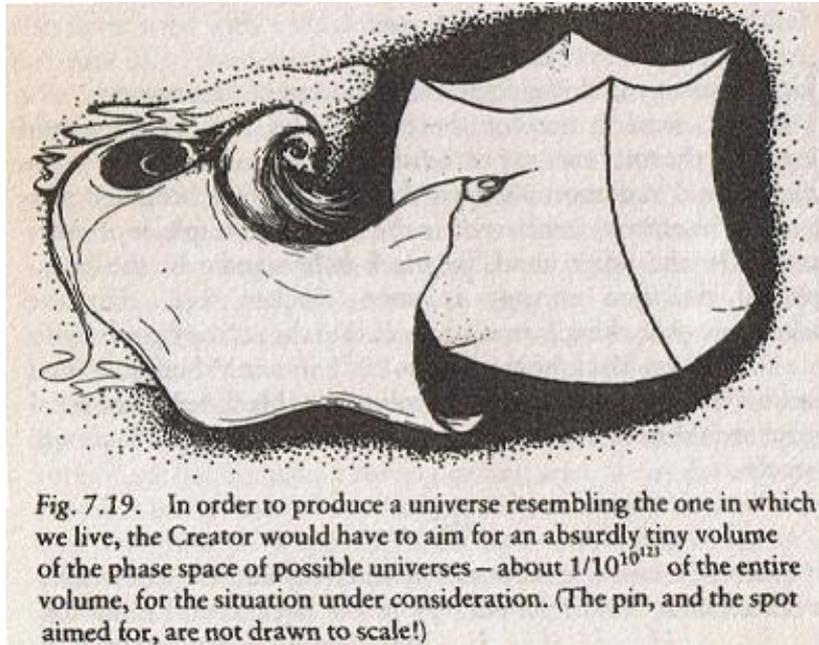
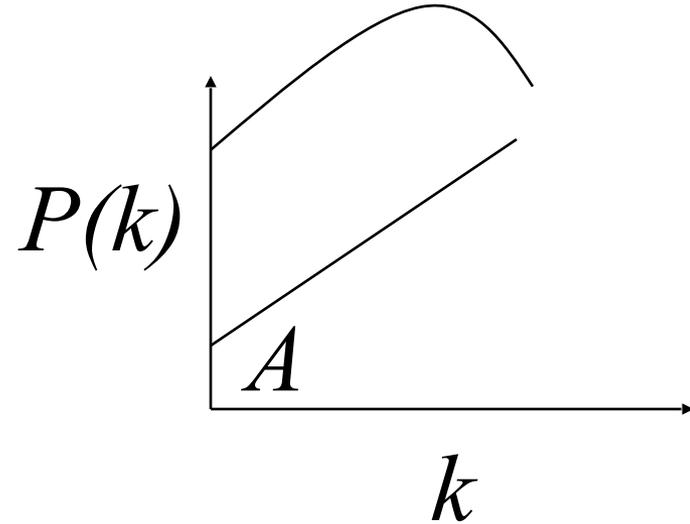


Fig. 7.19. In order to produce a universe resembling the one in which we live, the Creator would have to aim for an absurdly tiny volume of the phase space of possible universes – about $1/10^{10^{123}}$ of the entire volume, for the situation under consideration. (The pin, and the spot aimed for, are not drawn to scale!)

$$P(k, t) \sim A T^2(k, t) k^n$$



$$S(t) \sim P(t)?$$

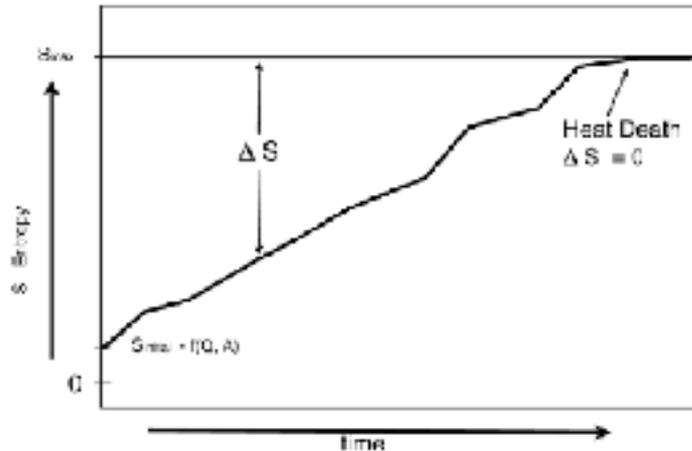
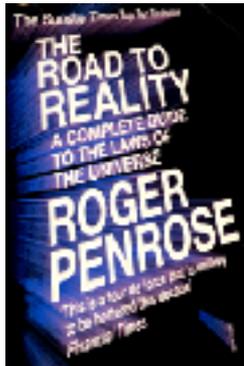


Fig. 9. The universe starts off at low entropy (not zero) due to the low level of density perturbations in the early universe—low Q and low A (e.g. [90])—where “low” means less than the maximum value S_{max} . At S_{max} all the energy density of the universe is in massless particles in equilibrium at a common temperature. Thus the universe starts off with a large entropy gap ΔS . The parameters Q and A are the observable normalizations of the primordial density fluctuations and set the initial gravitational entropy of the universe. There is no general agreement on the curve shown here. See for example Fig. 7.5 in [14] and Fig. 1.2 in [21].