The 3rd Korea-Japan Workshop on Dark Energy

# Redshift Space Distortion as a Standard Ruler

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# OUTLINE

- \* Power spectrum in redshift space
  - Fingers of God and Kaiser effect (Redshift space distortion)
- \* Monopole and Quadrupole in linear theory
- \* Redshift space perturbed Power spectrum
- \* Monopole and Quadrupole in quasi-linear theory
- \* Dark energy dependence on quadrupole

# POWER SPECTRUM IN REDSHIFT SPACE

- \* Kaiser effect (RSD): due to coherent peculiar velocities towards the central mass (large scale effect) (refer Matsubara's talk)
- \* Fingers of God effect: random velocity dispersions in galaxy clusters (affects only redshift not position, stretching only radially, small scale effect)

  Real space matter power spectrum

(what we calculate)

\* Redshift space linear power spectrum

$$P_s^{(L)}(k,\mu,z) = D_{FoG}^2(k,\mu,z)b(k,z)^2 (1 + \beta \mu^2)^2 P^{(L)}(k,\mu,z)$$

where 
$$\beta(k,z) = \frac{f(z)}{b(k,z)}$$
,  $f(z) = \frac{d \ln D_1}{d \ln a}$ ,  $\delta^{(1)}(k,z) = D_1(z)\delta_1(k)$ 

# RSD and Fingers of God

\* Redshift space distortion (RSD) is measured using (refer Okumura's talk)

where 
$$\beta(k,z) = \frac{f(z)}{b(k,z)}$$
,  $f(z) = \frac{d \ln D_1}{d \ln a}$ ,  $\delta^{(1)}(k,z) = D_1(z)\delta_1(k)$ 

Lorentzian and Gaussian forms

$$\begin{aligned} \mathrm{D}^{\mathrm{Lor}}_{\mathrm{FoG}}(k,\mu,z,\sigma) &=& \frac{1}{1+0.5[k\mu\sigma(z)D(z)f(z)]^2} \\ \mathrm{D}^{\mathrm{Gau}}_{\mathrm{FoG}}(k,\mu,z,\sigma) &=& \exp\left[-0.5[k\mu\sigma(z)D(z)f(z)]^2\right] \end{aligned}$$

velocity dispersion

$$\sigma^{2}(z) = \frac{4\pi}{3} \int \frac{dq}{(2\pi)^{3}} P_{\theta\theta}(q, z)$$

# Monopole and Quadrupole in Linear Theory

\* Decompose to angular dependence (multipoles)

$$P(k, \mu, z) = \sum_{l=0,2,4,\cdots} P_l(k, z) \mathcal{L}_l(\mu)$$
  $P_l(k, z) = \frac{2l+1}{2} \int_{-1}^1 d\mu P(k) \mathcal{L}_l(\mu)$ 

\* Monopole : angular averaged redshift space power spectrum

$$P_0(k,z) = (1 + \frac{2}{3}f + \frac{1}{5}f^2)P_L(k,z)$$
 Positive

\* Quadrupole : leading anisotropic contribution of redshift space power spectrum

$$P_2(k,z) = \left(\frac{4}{3}f + \frac{4}{7}f^2\right)P_L(k,z)$$
 Positive

# PERTURBED REDSHIFT SPACE POWER SPECTRUM

\* Perturbed redshift space power spectrum can be calculated from SPT/LPT/iPT (refer Matsubara's talk and Gong's talk for GR)

$$P_{s}^{(Scoccimarro)}(k,\mu,z) = \left(b^{2}(k,z)P_{\delta\delta}(k,z) + 2b(k,z)f(z)\mu^{2}P_{\delta\theta}(k,z) + f(z)^{2}\mu^{4}P_{\theta\theta}(k,z)\right)$$
$$= b^{2}(k)\left(P_{\delta\delta}(k) + 2\beta(k,z)\mu^{2}P_{\delta\theta}(k) + \beta^{2}(k,z)\mu^{4}P_{\theta\theta}\right) \qquad ($$

$$\begin{split} \mathbf{P}_{\mathrm{s}}^{\mathrm{(Matsubara)}}(k,\mu,z) &= \mathbf{P}_{\mathrm{s},11}(k,\mu,z) + \mathbf{P}_{\mathrm{s},22}(k,\mu,z) + 2\mathbf{P}_{\mathrm{s},13}(k,\mu,z) \\ &= \left[\mathbf{K}_{1}^{(\mathrm{s})}(k,\mu,z)\right]^{2} \mathbf{P}_{\mathrm{L}}(k,z) + 2\int \frac{d^{3}q}{(2\pi)^{3}} \mathbf{P}_{\mathrm{L}}(q,z) \mathbf{P}_{\mathrm{L}}(|k-q|,z) \left[\mathbf{K}_{2}^{(\mathrm{s})}(\vec{q},\vec{k}-\vec{q},\mu,z)\right]^{2} \\ &+ 6\left[\mathbf{K}_{1}^{(\mathrm{s})}(k)\right] \mathbf{P}_{\mathrm{L}}(k,z) \int \frac{d^{3}q}{(2\pi)^{3}} \mathbf{P}_{\mathrm{L}}(q,z) \left[\mathbf{K}_{3}^{(\mathrm{s})}(\vec{q},-\vec{q},\vec{k},\mu,z)\right] \end{split} \tag{2}$$

$$P_{s}^{(TNS)}(k,\mu,z) = P_{s}^{(Scoccimarro)}(k,\mu,z) + A(k,\mu,b,z) + B(k,\mu,b,z)$$

Taruya, Nichimichi, Saito

# Monopole and Quadrupole in Quasi-linear Theory

#### \* Monopole

$$P_{0}(k,z) = \left(P_{L} + A_{00}P_{22} + B_{00}P_{13}\right) + \left(\frac{2}{3}P_{L} + \frac{1}{3}A_{11}P_{22} + \frac{1}{3}(B_{00} + B_{11})P_{13}\right)f$$

$$+ \left(\frac{1}{5}P_{L} + \left(\frac{1}{3}A_{12} + \frac{1}{5}A_{22}\right)P_{22} + \left(\frac{1}{3}B_{12} + \frac{1}{5}(B_{11} + B_{22})\right)P_{13}\right)f^{2}$$

$$+ \left(\left(\frac{1}{5}A_{23} + \frac{1}{7}A_{33}\right)P_{22} + \left(\frac{1}{7}B_{22} + \frac{1}{5}(B_{12} + B_{23})\right)P_{13}\right)f^{3}$$

$$+ \left(\left(\frac{1}{5}A_{24} + \frac{1}{7}A_{34} + \frac{1}{9}A_{44}\right)P_{22} + \frac{1}{7}B_{23}P_{13}\right)f^{4}$$

Positive

# Monopole and Quadrupole in Quasi-linear Theory

#### \* Quadrupole

$$P_{2}(k,z) = \left(\frac{4}{3}P_{L} + \frac{2}{3}A_{11}P_{22} + \frac{2}{3}(B_{00} + B_{11})P_{13}\right)f$$

$$+ \left(\frac{4}{7}P_{L} + \left(\frac{2}{3}A_{12} + \frac{4}{7}A_{22}\right)P_{22} + \left(\frac{4}{7}B_{11} + \frac{2}{3}B_{12} + \frac{4}{7}B_{22}\right)P_{13}\right)f^{2}$$

$$+ \left(\left(\frac{4}{7}A_{23} + \frac{10}{21}A_{33}\right)P_{22} + \left(\frac{4}{7}B_{12} + \frac{10}{21}B_{22} + \frac{4}{7}B_{23}\right)P_{13}\right)f^{3}$$

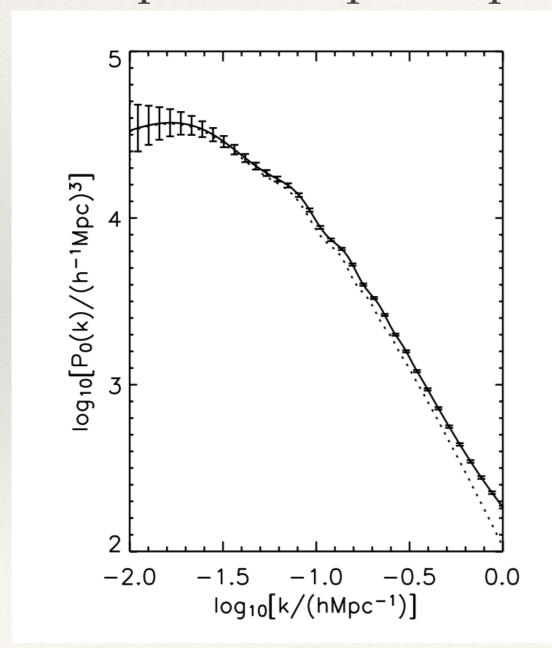
$$+ \left(\left(\frac{4}{7}A_{24} + \frac{10}{21}A_{34} + \frac{40}{99}A_{44}\right)P_{22} + \frac{10}{21}B_{23}P_{13}\right)f^{4}$$

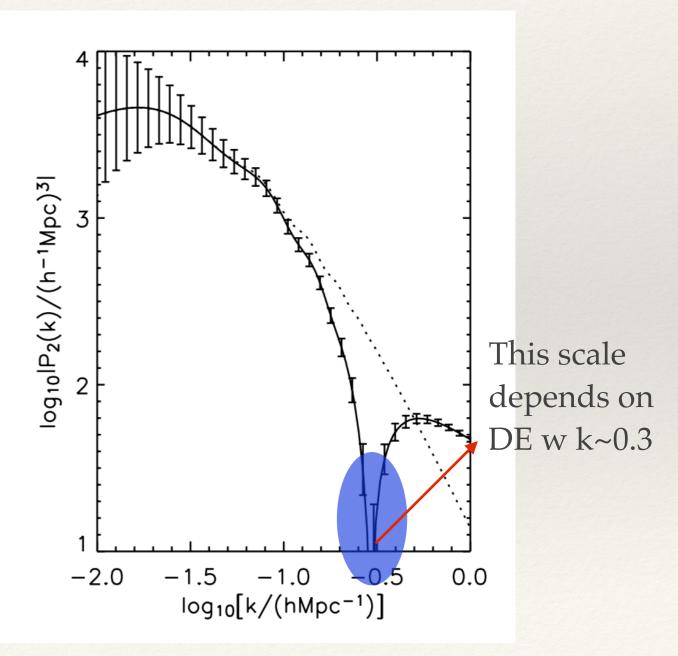
This term change signs at the specific scale k for the different DE models

Positive → Negative as k ↑

# Monopole and Quadrupole in Quasi-linear Theory

\* Monopole and quadrupole Yamamoto, Bassett, Nishioka 05





# Dark Energy Dependence on Quadrupole

$$\begin{split} \mathbf{P}_{\mathrm{s}}^{\mathrm{(Matsubara)}}(k,\mu,z) &= \mathbf{P}_{\mathrm{s},11}(k,\mu,z) + \mathbf{P}_{\mathrm{s},22}(k,\mu,z) + 2\mathbf{P}_{\mathrm{s},13}(k,\mu,z) \\ &= \left[ \mathbf{K}_{1}^{(\mathrm{s})}(k,\mu,z) \right]^{2} \mathbf{P}_{\mathrm{L}}(k,z) + 2 \int \frac{d^{3}q}{(2\pi)^{3}} \mathbf{P}_{\mathrm{L}}(q,z) \mathbf{P}_{\mathrm{L}}(|k-q|,z) \left[ \mathbf{K}_{2}^{(\mathrm{s})}(\vec{q},\vec{k}-\vec{q},\mu,z) \right]^{2} \\ &+ 6 \left[ \mathbf{K}_{1}^{(\mathrm{s})}(k) \right] \mathbf{P}_{\mathrm{L}}(k,z) \int \frac{d^{3}q}{(2\pi)^{3}} \mathbf{P}_{\mathrm{L}}(q,z) \left[ \mathbf{K}_{3}^{(\mathrm{s})}(\vec{q},-\vec{q},\vec{k},\mu,z) \right] \end{split} \tag{2}$$

\* 
$$P_{11}$$
:  $K_1^{(s)}(k,\mu,z) = 1 + \frac{f(w,z)}{b}\mu^2$   $P_L(k,z) = D_1(w,z)^2 P_L(k,w)$ 

$$F_2(a, \vec{q}, \vec{k} - \vec{q}) = \frac{2c_{22}(r - x) + c_{21}(r + x - 2rx^2)}{2r(1 + r^2 - 2rx)}$$

$$G_2(a, \vec{q}, \vec{k} - \vec{q}) = -\frac{2c_{\theta 22}(r - x) + c_{\theta 21}(r + x - 2rx^2)}{2r(1 + r^2 - 2rx)}$$

Ex: 
$$\frac{A_{00}}{(1+r^2-2rx)^2} = 2r^2F_2^2 = \frac{(7x+r(3-10x^2))^2}{98(1+r^2-2rx)^2}$$

DE dependence on  $c_{21}(z,w)$ ,  $c_{22}(z,w)$ , etc Usually, one considers w dependence on  $D_1$  only

## Improvement by Including T-dependence on F & G

\* Comparison on the matter power spectrums with and without EdS assumption in the kernels (SL, Park, Biern 14)

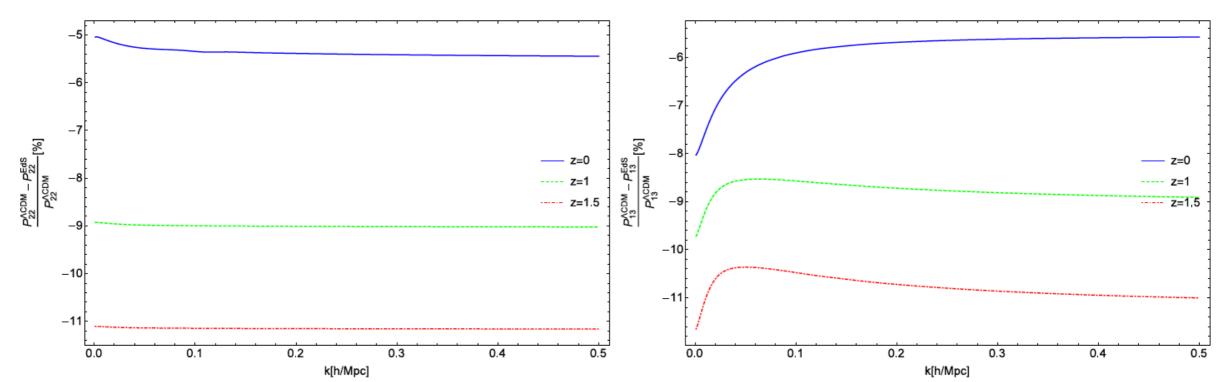


FIG. 2: Errors in  $P_{22}$  and  $P_{13}$  a) Differences between the correct  $P_{22}$  and the one with EdS assumption at the different epoches. The solid, dashed, and dotdashed lines correspond to z = 0, 1.0, and 1.5, respectively. a) Differences between the correct  $P_{13}$  and the EdS assumed  $P_{13}$  at different epoches.

### Improvement by Including T-dependence on F & G

 Comparison on the matter power spectrums with and without EdS assumption in the kernels

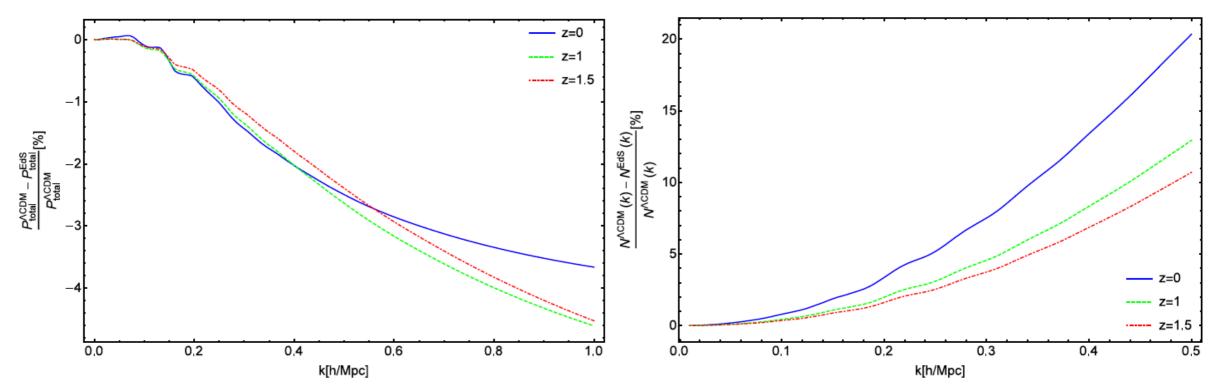
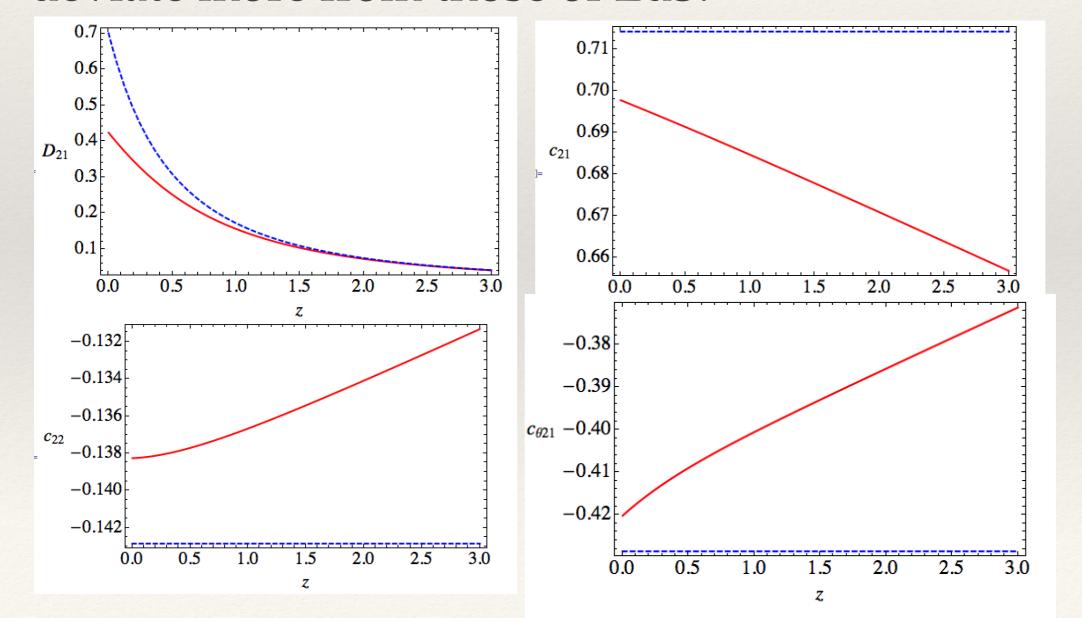


FIG. 3: Errors in  $P_{\text{total}}$  and N a) Differences between the correct  $P_{\text{total}}$  and the one with  $\lambda = 1$  (EdS) assumption at the different epoches. The solid, dashed, and dotdashed lines correspond to z = 0, 0.5, and 1, respectively. a) Differences between the correct  $P_{NL}$  and the  $\lambda = 1$  (EdS) assumed  $P_{NL}$  at different epoches.

# Improvement by Including T-dependence on F & G

\* As z increases, D approaches to that of EdS. But F & G deviate more from those of EdS.



## CONTINUATION

- \* How sensitive the flipping scale on DE e.o.s
- \* Comparison with other sensitivities ( $\Omega_m$ , b, etc)
- \* RS effect has similar properties but with very low S/N (*SL. 15*)

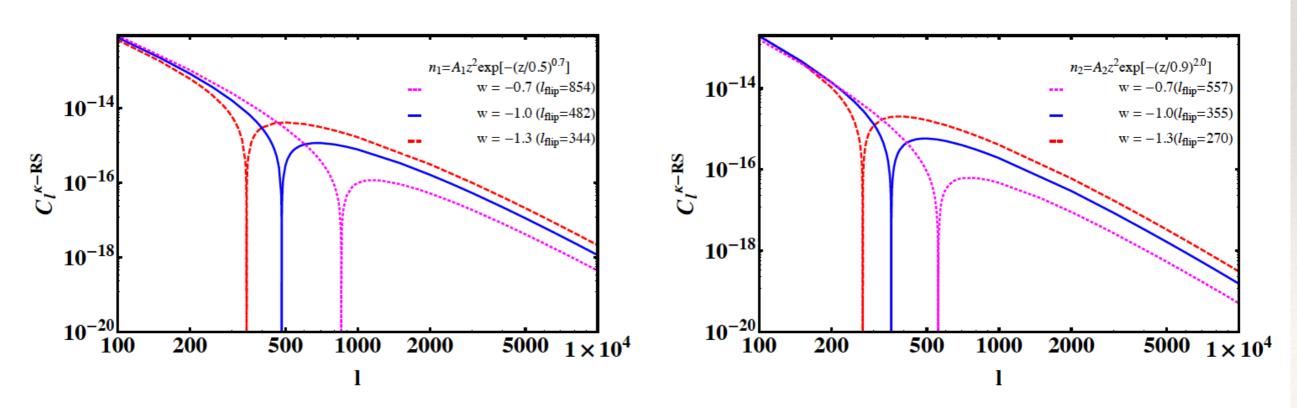


FIG. 1: CMB-WL cross-correlation spectra,  $C_{i}^{RS-\kappa}$  of different dark energy models for different WL surveys  $(n_{1(2)})$ .

7##MXS! ありがとうございます! 감사합니다!