

# Study on the mapping of dark matter clustering from real space to redshift space

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Yi Zheng, Yong-Seon Song, arXiv:1603.00101

### **RSD** Introduction

Redshift space distortion

In observation, galaxy distance is determined by "redshift"

**Peculiar velocity** of galaxies cause them to appear displaced along the line of sight in redshift space

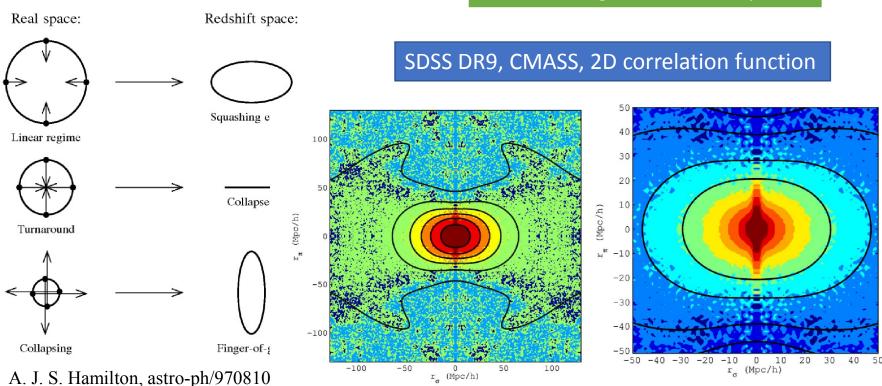
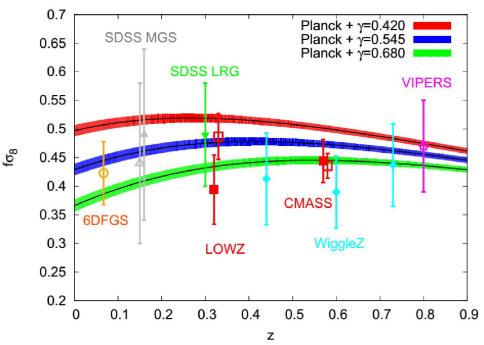


Figure 3. Left panel: Two-dimensional correlation function of CMASS galaxies (color) compared with the best fit model described in Section [5.1] (black lines). Contours of equal  $\xi$  are shown at [0.6, 0.2, 0.1, 0.05, 0.02, 0]. Right panel: Smaller-scale two-dimensional clustering. We show model contours at [0.14, 0.05, 0.01, 0]. The value of  $\xi_0$  at the minimum separation bin in our analysis is shown as the innermost contour. The  $\mu \approx 1$  "finger-of-god" effects are small on the scales we use in this analysis.

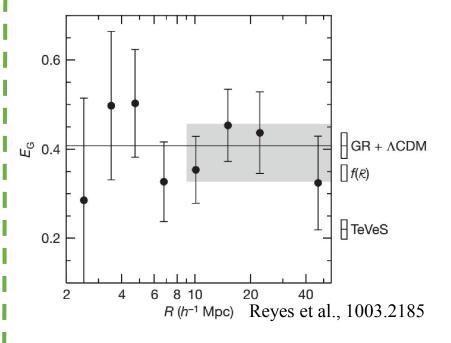
## **RSD** Introduction

$$f = \frac{d \ln D}{d \ln a} = \Omega_m(z)^{\gamma}$$
 Eric Linder, astro-ph/0507263



Gil-Marin et al. 1509.06386

$$\langle \hat{E}_G \rangle = \left[ \frac{\nabla^2 (\psi - \phi)}{-3H_0^2 a^{-1} \theta} \right]_{k=l/\bar{\chi},\bar{z}} = \left[ \frac{\nabla^2 (\psi - \phi)}{3H_0^2 a^{-1} \beta \delta} \right]_{k=l/\bar{\chi},\bar{z}} \equiv E_G.$$
Zhang et al., 0704.1932



#### **RSD** correlation function

#### Streaming model:

$$1 + \xi_S(s_\perp, s_\parallel) = \int dr_\parallel \left[ 1 + \xi_R(r) \right] \mathcal{P}(r_\parallel - s_\parallel | \mathbf{r})$$

Improving the modelling of redshift-space distortions:

I. A bivariate Gaussian description for the galaxy pairwise velocity distributions

arXiv:1407.4753v2

Davide Bianchi<sup>1,2\*</sup>, Matteo Chiesa<sup>1,2</sup> & Luigi Guzzo<sup>1</sup>

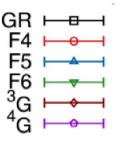
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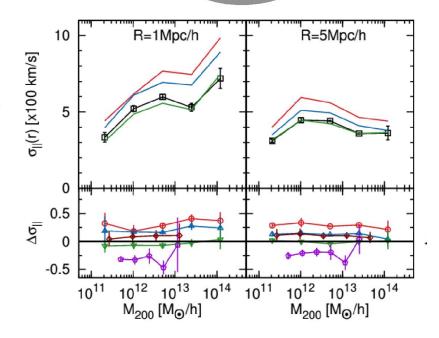
Clear and Measurable Signature of Modified Gravity in the Galaxy

Velocity Field

Wojciech A. Hellwing, Alexandre Barreira, Carlos S. Frenk, Baojiu Li, and Shaun Cole Phys. Rev. Lett. 112, 221102 - Published 5 June 2014



Pairwise velocity probability distribution function



## Power spectrum modelling: 1. mapping; 2. PT

We follow the derivation of TNS paper, Taruya et al. 1006.0699

The redshift position **s**:

$$\mathbf{s} = \mathbf{x} + v_{z}/(aH)\,\hat{\mathbf{z}}$$

From mass conservation:

$$(1 + \delta^s(\mathbf{s}))d^3s =$$

Taylor expansion and trunction

The redshift space power spectrum:

$$P^{(S)}(\mathbf{k}) = \int d^3 \mathbf{x} e^{i\mathbf{k}\cdot\mathbf{x}} \langle e^{-ik\mu f\Delta u_z} \{ \delta(\mathbf{r}) + f\nabla_z u_z(\mathbf{r}) \}$$

$$\times \{ \delta(\mathbf{r}') + f\nabla_z u_z(\mathbf{r}') \} \rangle,$$

$$u_z(\mathbf{r}) = -v_z(\mathbf{r})/(aHf)$$

## **RSD** modelling

$$P^{(S)}(\mathbf{k}) = \int d^3 \mathbf{x} e^{i\mathbf{k}\cdot\mathbf{x}} \langle e^{-ik\mu f\Delta u_z} \{ \delta(\mathbf{r}) + f\nabla_z u_z(\mathbf{r}) \}$$
$$\times \{ \delta(\mathbf{r}') + f\nabla_z u_z(\mathbf{r}') \} \rangle,$$

#### Definition:

$$j_1 = -ik\mu f, \qquad A_1 = u_z(\mathbf{r}) - u_z(\mathbf{r}'),$$

$$A_2 = \delta(\mathbf{r}) + f\nabla_z u_z(\mathbf{r}), \qquad A_3 = \delta(\mathbf{r}') + f\nabla_z u_z(\mathbf{r}').$$

$$P^{(S)}(k,\mu) = \int d^3x e^{ik \cdot x} \langle e^{j_1 A_1} A_2 A_3 \rangle,$$

$$\begin{split} \langle e^{j_1 A_1} A_2 A_3 \rangle &= \exp\{\langle e^{j_1 A_1} \rangle_c\} [\langle e^{j_1 A_1} A_2 A_3 \rangle_c \\ &+ \langle e^{j_1 A_1} A_2 \rangle_c \langle e^{j_1 A_1} A_3 \rangle_c]. \end{split}$$



Important for separation of FoG

$$P^{(S)}(k,\mu) = \int d^3x e^{ik\cdot x} \exp\{\langle e^{j_1 A_1} \rangle_c\} [\langle e^{j_1 A_1} A_2 A_3 \rangle_c] + \langle e^{j_1 A_1} A_2 \rangle_c \langle e^{j_1 A_1} A_3 \rangle_c].$$

Scale dependent part+scale independent part

$$\begin{split} \langle e^{j_1 A_1} A_2 A_3 \rangle_c &+ \langle e^{j_1 A_1} A_2 \rangle_c \langle e^{j_1 A_1} A_3 \rangle_c \\ &\simeq \langle A_2 A_3 \rangle + j_1 \langle A_1 A_2 A_3 \rangle_c \\ &+ j_1^2 \{ \frac{1}{2} \langle A_1^2 A_2 A_3 \rangle_c + \langle A_1 A_2 \rangle_c \langle A_1 A_3 \rangle_c \} + \mathcal{O}(j_1^3). \end{split}$$

#### Final model

#### 1006.0699 & 1603.00101

Taylor expansion in terms of  $j_1 = -ik\mu$ 

$$D_{\text{local}}^{\text{FoG}}(k\mu, \boldsymbol{x}) \left[ \langle e^{j_1 A_1} A_2 A_3 \rangle_c + \langle e^{j_1 A_1} A_2 \rangle_c \langle e^{j_1 A_1} A_3 \rangle_c \right]$$

$$\simeq j_1^0 \langle A_2 A_3 \rangle_c + j_1^1 \langle A_1 A_2 A_3 \rangle_c$$

$$+ j_1^2 \left\{ \langle A_1 A_2 \rangle_c \langle A_1 A_3 \rangle_c + \frac{1}{2} \langle A_1^2 A_2 A_3 \rangle_c - \langle u_z u_z' \rangle_c \langle A_2 A_3 \rangle_c \right\}$$

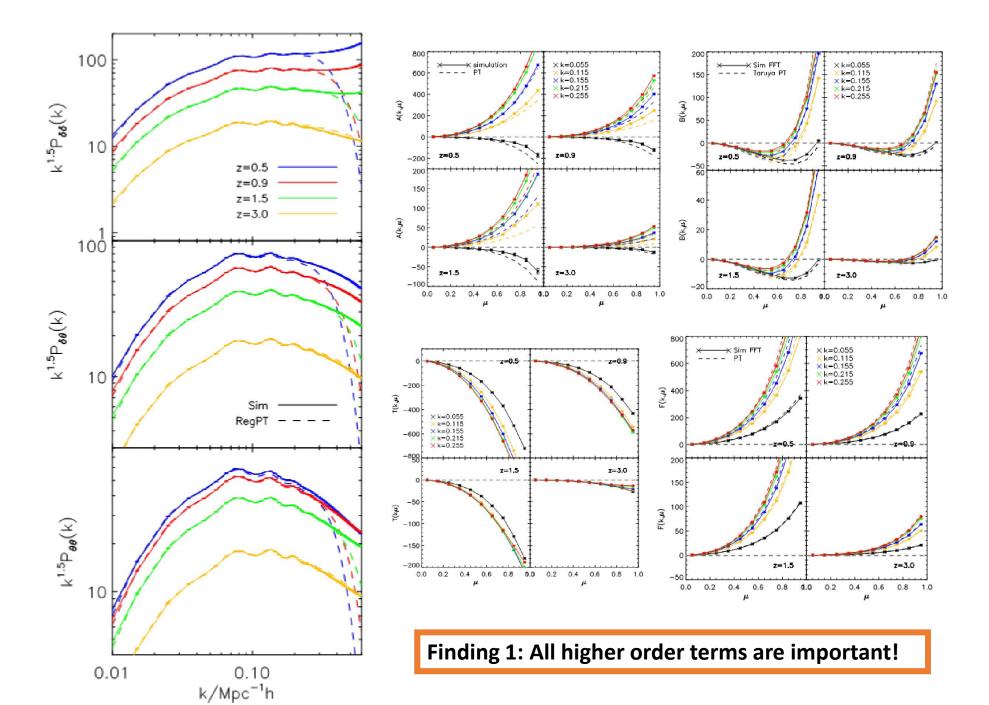
$$+ \mathcal{O}(j_1^3), \qquad (14)$$

$$P^{(S)}(k,\mu) = D^{FoG}(k\mu\sigma_z)P_{perturbed}(k,\mu)$$

$$= D^{FoG}(k\mu\sigma_z)[P_{\delta\delta} + 2\mu^2 P_{\delta\Theta} + \mu^4 P_{\Theta\Theta} (15)$$

$$+A(k,\mu) + B(k,\mu) + T(k,\mu) + F(k,\mu)].$$

parameter	physical meaning	value
$\Omega_m$	present fractional matter density	0.3132
$\Omega_{\Lambda}$	$1 - \Omega_m$	0.6868
$\Omega_b$	present fractional baryon density	0.049
h	$H_0/(100 \text{ km s}^{-1}\text{Mpc}^{-1})$	0.6731
$n_s$	primordial power spectral index	0.9655
$\sigma_8$	r.m.s. linear density fluctuation	0.829
$L_{\rm box}$	simulation box size	$1890 \ h^{-1}{\rm Mpc}$
$N_{ m p}$	simulation particle number	$1024^3$
$m_{ m p}$	simulation particle mass	$5.46 \times 10^{10} h^{-1} M_{\odot}$
NT.	1	10

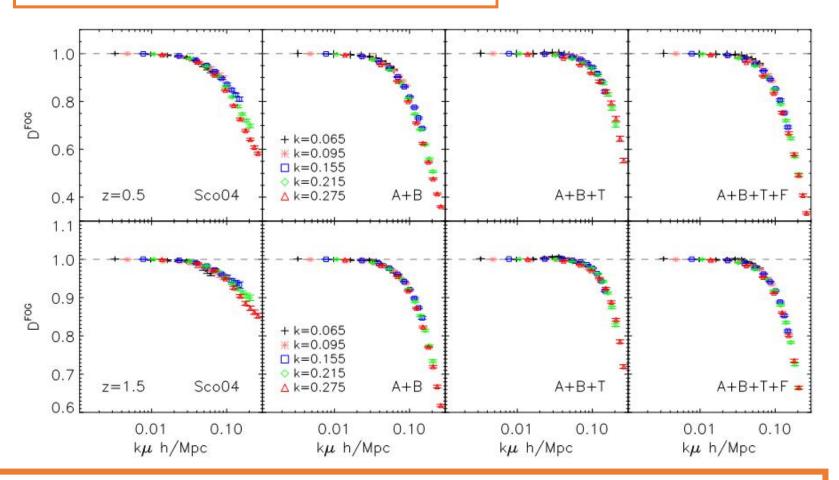


### 1. FoG test

$$P^{(S)}(k,\mu) = D^{FoG}(k\mu\sigma_z)P_{perturbed}(k,\mu)$$
  
=  $D^{FoG}(k\mu\sigma_z)[P_{\delta\delta} + 2\mu^2P_{\delta\Theta} + \mu^4P_{\Theta\Theta}$  (15)  
+  $A(k,\mu) + B(k,\mu) + T(k,\mu) + F(k,\mu)].$ 

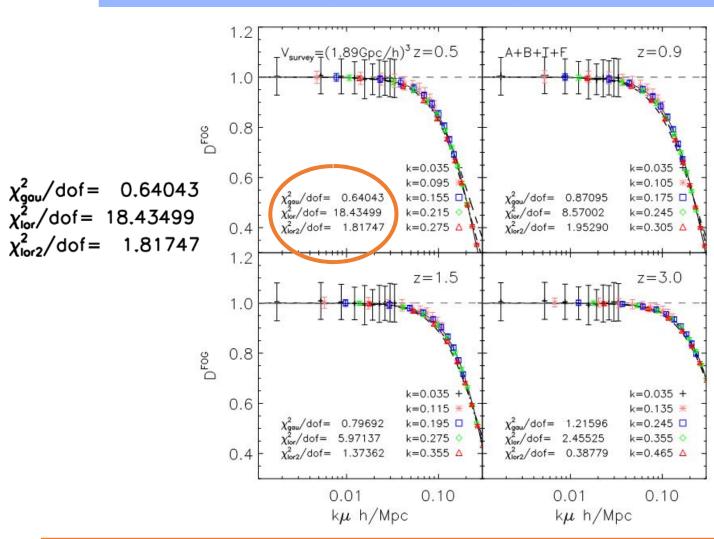
One 'standard' for good RSD model: if FoG term is a simple function of  $k\mu\sigma_z$ 

$$D^{\text{FoG}} = \frac{P_{\text{sim}}^{(S)}(k,\mu)}{P_{\text{perturbed}}(k,\mu)},$$



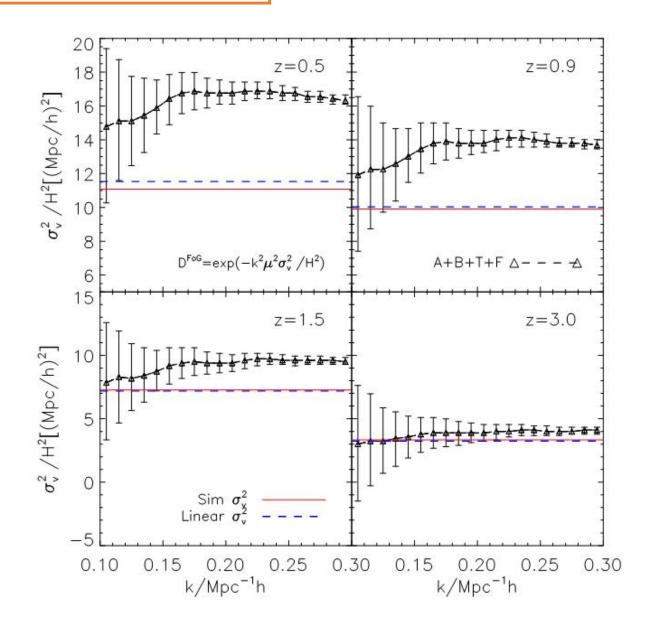
Finding 2: A+B+F+T works best, also in order to be complete at j\_1^2 order, we will consider A+B+F+T model hereafter.

$$D^{\text{FoG}}(k\mu\sigma_z) = \begin{cases} e^{-k^2\mu^2\sigma_z^2/2} & \text{Gaussian,} \\ \left(1 + k^2\mu^2\sigma_z^2/2\right)^{-1} & \text{Lorentzian,} \\ \left(1 + k^2\mu^2\sigma_z^2/2\right)^{-2} & \text{squared Lorentzian,} \end{cases}$$



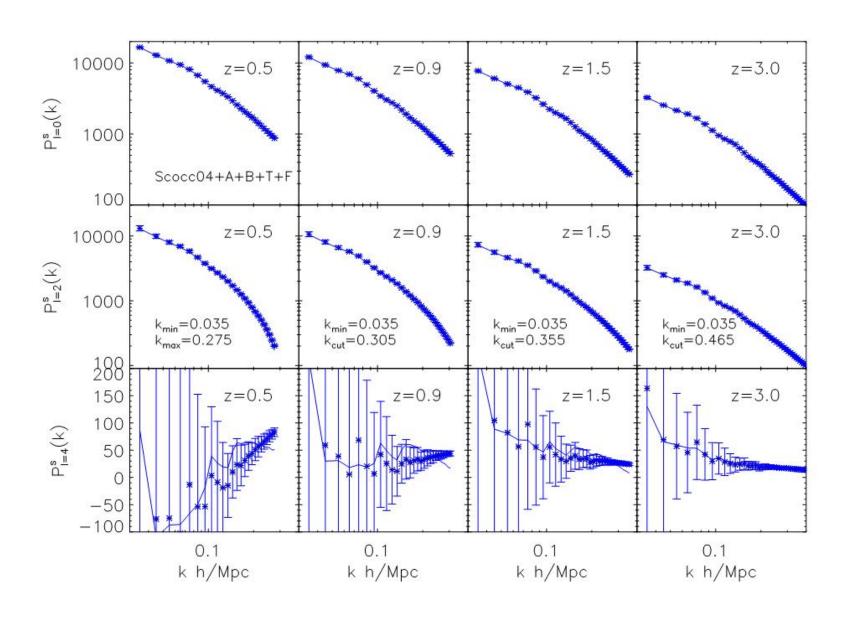
Finding 3: Gaussian model is best FoG model for A+B+F+T case

#### Check if $\sigma_z$ is a constant over scales:

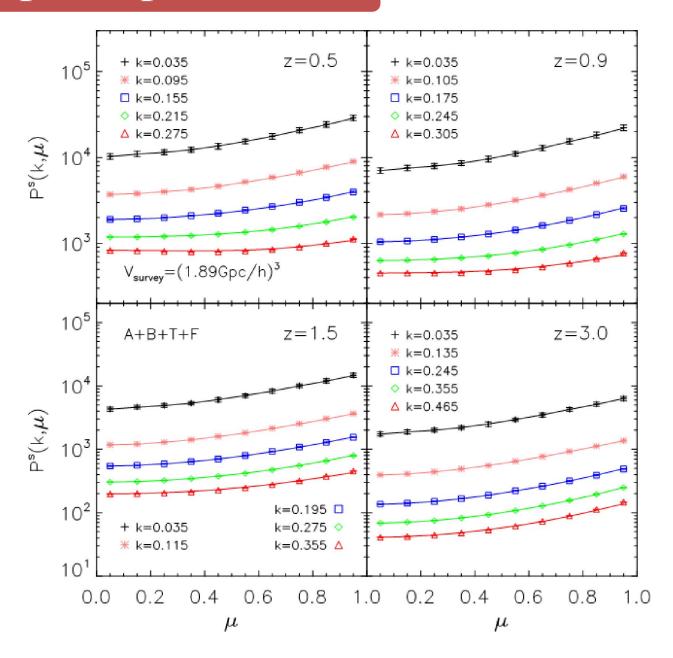


# 2. Multipole test:

$$P_{\ell}^{(S)}(k) = \frac{2\ell+1}{2} \int_{-1}^{1} d\mu P^{(S)}(k,\mu) \mathcal{P}_{\ell}(\mu),$$

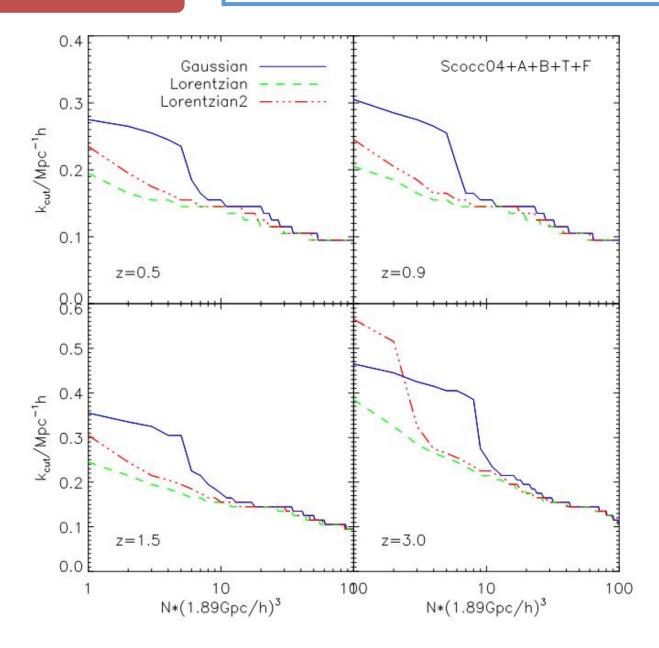


## 3. 2D power spectrum test:



# 4. k\_cut test:

The scale up to which systemic error = statistical error



## Future work

- Test the mapping for halo case and galaxy case
- Develop robust halo/galaxy bias scheme, so combine the bias model with dark matter template measured from simulations, we could do the cosmological constrain from data.

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